

# SUPERFLUID $^3\text{He}$ : SOME PRE-HISTORY

$^4\text{He}$ : below 2K, **superfluid**

## ELECTRONS IN METALS

since ancient times

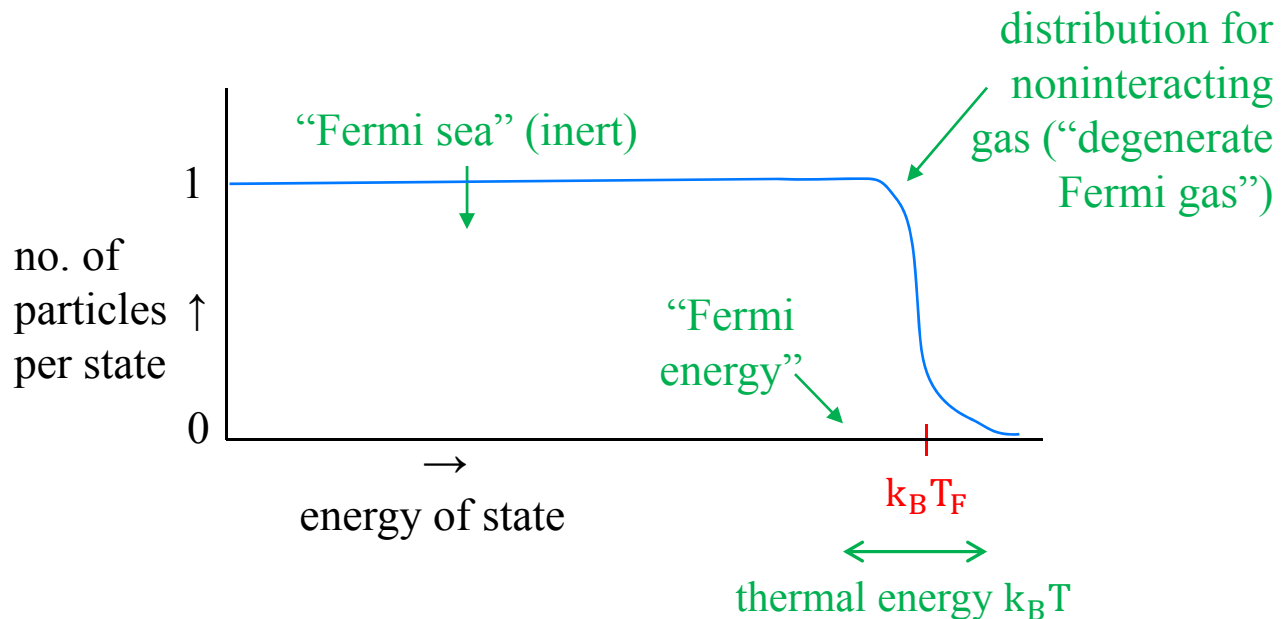
charged

## LIQUID $^3\text{He}$

since ~ 1950

neutral

Particles of spin  $\frac{1}{2} \Rightarrow$  **Fermi statistics**



Landau (1957): interactions don't change picture qualitatively (in "normal" phase) ("degenerate Fermi liquid")

$T_F \sim 10^4 - 10^5 \text{ K}$

at  $T \lesssim 20\text{K}$ .

**superconductivity**

(in some metals)

$T_F \sim 5 \text{ K}$

at  $T \lesssim 10^{-3}\text{K}$ .

**superfluidity??**

## MORE PREHISTORY

Theory of superconductivity:

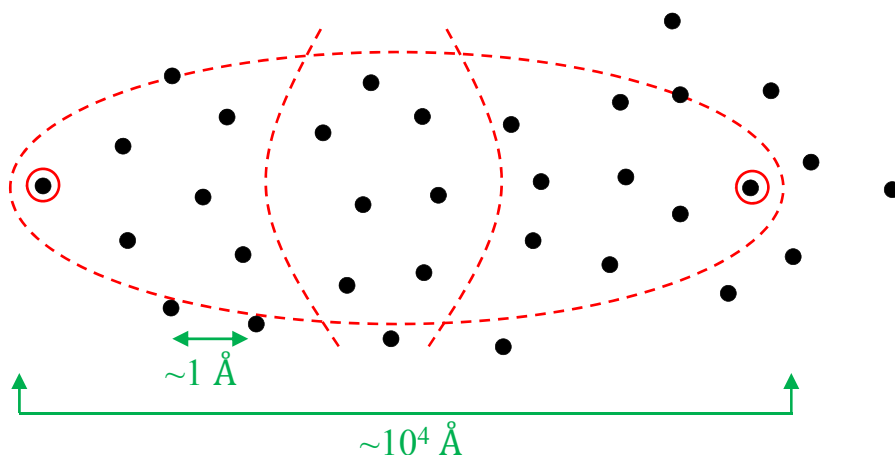
- a) phenomenological (V. L. Ginzburg, A. A. Abrikosov, et al., (1950-1955):

macroscopic wave function

- b) microscopic (Bardeen et al., 1957):

electrons in energy shell of width  $\sim k_B T_C$  around Fermi energy form Cooper pairs

critical temp.,  $\lesssim 20\text{K}$



Crucial feature of BCS theory: ALL COOPER PAIRS MUST BEHAVE IN EXACTLY THE SAME WAY!

(GL “macroscopic wave function” is just the common center-of-mass wave function of all the pairs)

In BCS theory, “internal” wave function of pairs trivial: (“ $^1S_0$ ”)

$$\psi(\underline{r}_1 \underline{r}_2; \sigma_1 \sigma_2) \sim \frac{1}{\sqrt{2}} (\uparrow_1 \downarrow_2 - \downarrow_1 \uparrow_2) f(|\underline{r}_1 - \underline{r}_2|)$$

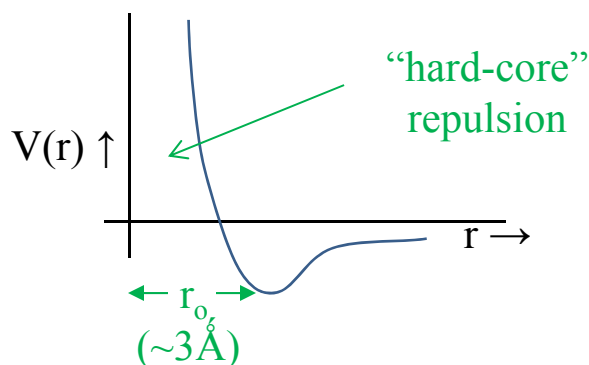
spin singlet

spherically symmetric

( $\ell = 0$ )

NO INTERNAL (“ORIENTATIONAL”) DEGREES OF FREEDOM

## EARLY THEORETICAL WORK ON POSSIBLE COOPER PAIRING IN LIQUID $^3\text{He}$



$$r \sim r_0, p \sim p_F (\equiv \sqrt{2mk_B T_F})$$

$$\Rightarrow \text{relative angular momentum}$$

$$\ell \equiv (p_F r_0 / \hbar) \neq 0$$

$$(\text{prob. 1 or 2})$$

Pauli principle:  $\begin{cases} \ell = 0, 2, 4 \dots & S = 0 \text{ (singlet)} \\ \ell = 1, 3, 5 \dots & S = 1 \text{ (triplet)} \end{cases}$

in general,  $\ell \neq 0 \Rightarrow$  relative (internal) wave function of pair “equal spin pairing”  
has orientational degree(s) of freedom!

Anderson & Morel (1961): explore in detail case  $\ell = 2$ , and a special case of  $\ell = 1$ : only  $\uparrow\uparrow$  and  $\downarrow\downarrow$  pairs form, and have the same orbital angular momentum in direction  $\hat{\ell}$  (“ABM” state) Physical properties anisotropic.

Vdovin  
 Balian & Werthamer } (1963): in  $\ell = 1$  case all spin components “ $^3P_0$ ”

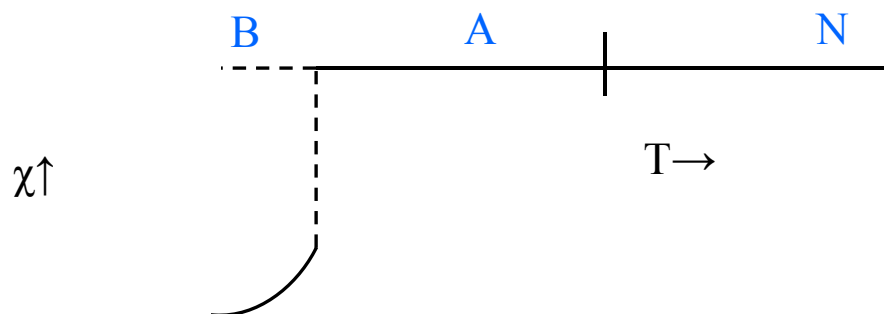
$(\uparrow\uparrow, \downarrow\downarrow, \frac{1}{\sqrt{2}} \uparrow\downarrow + \downarrow\uparrow)$  can form: in fact for any given pair,  $\underline{L} = -\underline{S} \Rightarrow J = 0$ .

(“BW” state). All physical properties isotropic. More stable than any ESP state.

Theoretical expectation c. 1964:

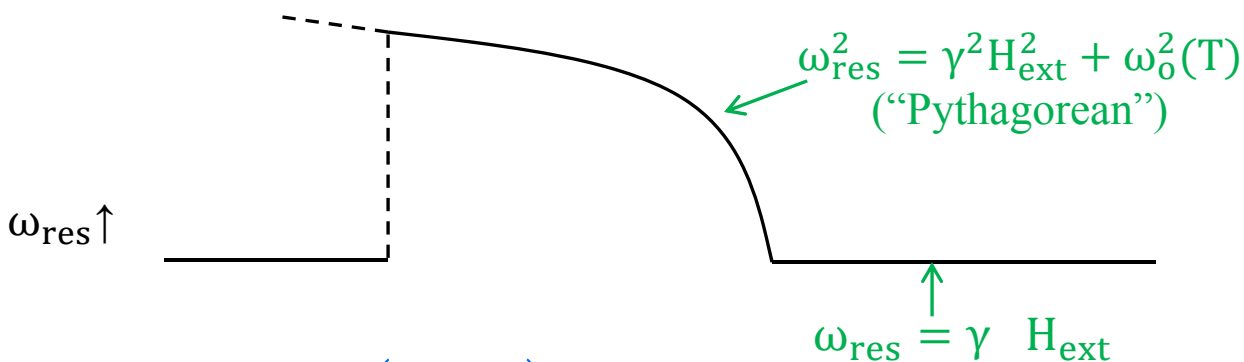
Liquid  $^3\text{He}$  may form Cooper pairs, either  $\ell = \text{even}$  (spin singlet) or with  $\ell = \text{odd}$  (BW state). In either case,  $\chi$  reduced and all magnetic properties isotropic.  $T_c$  difficult to predict.

NMR in the new phases:

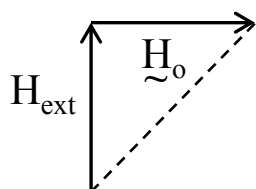


Not necessarily mysterious: e.g. A phase could be an ESP state (only  $\uparrow\uparrow, \downarrow\downarrow$  pairs  $\Rightarrow$  no reduction in  $\chi$ ), B could be singlet or BW (some  $\uparrow\downarrow$  pairs, so  $\chi$  reduced) [but: why is ESP ever stable?]-

But: what about the resonance frequency?



$$\omega_0^2(T) \approx A \left(1 - \frac{T}{T_A}\right), \quad \frac{A}{(2\pi)^2} \cong 5 \times 10^{10} \text{Hz}^2$$



$(\equiv \omega_0(T)/\gamma)$   
 Need  $H_0 \sim 30\text{G}$ . But, only spin-nonconserving force in problem is **nuclear dipole-dipole interaction**, and max. associated field is  $< 1\text{G}$ !

**IS THIS THE FIRST INDICATION OF A RADICAL BREAKDOWN OF QUANTUM MECHANICS?**

## WHAT CAN BE INFERRED FROM SUM RULES?

**IF** a single sharp resonance is observed (as in expt.) then:

$$\omega_{\text{res}}^2 = \gamma^2 H_{\text{ext}}^2 + \omega_0^2$$

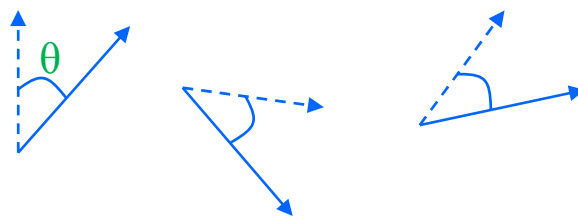
$$\omega_0^2 = \frac{\gamma^2 \chi^{-1} \partial^2 \langle H_D \rangle}{\partial \theta^2}$$

But  $\frac{\partial^2 \langle H_D \rangle}{\partial \theta^2} \sim \langle H_D \rangle$ :

So, exptl. value of  $\omega_0^2(T) \Rightarrow$

$$\langle H_D \rangle(T) \sim K \left( 1 - \frac{T}{T_A} \right), K \sim 10^{-3} \frac{\text{ergs}}{\text{cm}^3}$$

nuclear dipole energy



angle of simultaneous rot.<sup>n</sup> of all spins

HOW CAN THIS BE?

{ ↑ (“bad”) ↑  
→ (“good”) →

$$\Delta E \lesssim \frac{\mu_0 \mu_n^2}{r_0^3} \sim 10^{-7} \text{ K} \ll k_B T$$

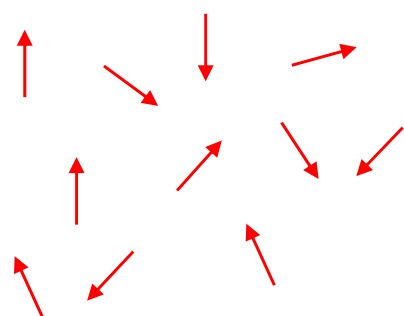
So, prima facie, preference for “good” orientation over “bad” is

at most  $\sim \frac{\Delta E}{k_B T} \sim 10^{-4}$  [actually,  $\sim \frac{\Delta E}{k_B T_F} \sim 10^{-7}$ ]

$\Rightarrow$  expectation value of dipole energy much too small!

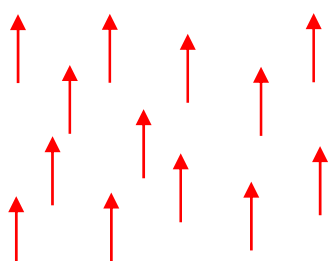
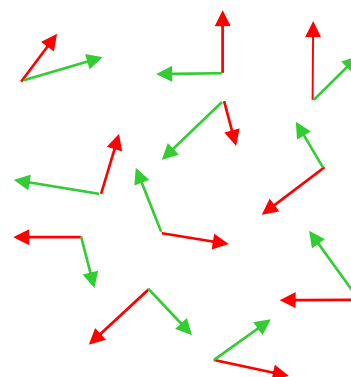
## SBSOS: ORDERING MAY BE SUBTLE

### FERROMAGNET

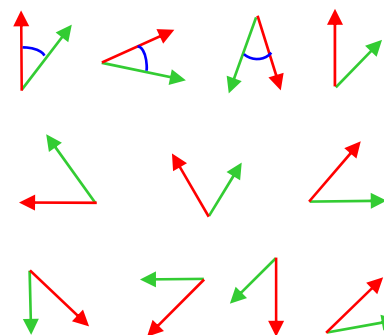


← NORMAL  
PHASE →

### LIQUID $^3\text{He}$



← ORDERED  
PHASE →



= total spin of pair  
 = relative orbital  
 ang. momentum)

$$\langle \tilde{S} \rangle \neq 0$$

$$\langle \tilde{S} \rangle = \langle \tilde{L} \rangle = 0$$

$$\text{but } \langle \tilde{L} \times \tilde{S} \rangle \neq 0!$$

“Absolute” direction  
chosen by (ultraweak)  
magnetic field

Relative direction chosen by  
(ultraweak) nuclear dipole  
force

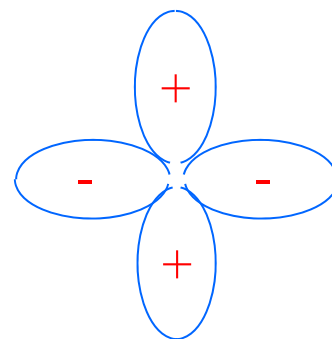
# MORAL OF $^3\text{He}$ NMR THEORY

“Behavioral conformism” of Cooper pairs applies not just to COM behavior but to **relative** motion of the components of a pair

⇒ in systems with “exotic” (non-s-wave) OP’s, spectacular **amplification** of the effects of ultra-weak forces (**dipole**, ? magnetic, ? PNC)

Other systems with “exotic” Cooper pairing:

U Pt <sub>3</sub> (etc.)	odd-parity
Sr <sub>2</sub> RuO <sub>4</sub>	p + ip (?)
ferropnictides	s <sub>±</sub> (?)
<b>cuprates</b>	<b>d<sub>x<sup>2</sup>-y<sup>2</sup></sub></b>



most “spectacular” amplification effects in cuprates:

Josephson effect

