

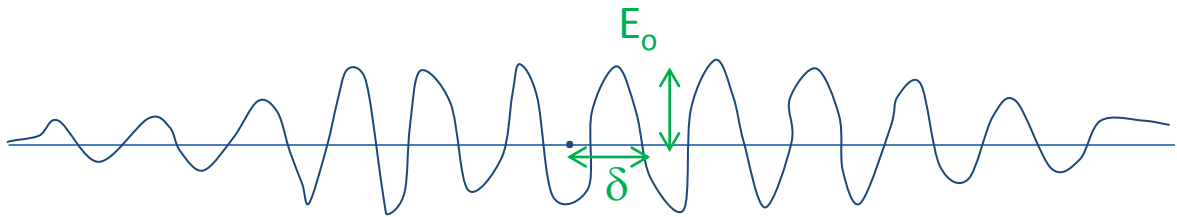
# “BASIC QUANTUM MECHANICS, AND SOME SURPRISING CONSEQUENCES”

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## CLASSICAL LIGHT WAVES

A (classical) light beam is a wave in which the electric field oscillates:

if I could “photograph” the electric field at a given time, it might look something like



and if I sit at a given point in space, it would have the same behavior in time: formally,

$$E(t) \cong E_0 \cos(\omega t + \delta)$$

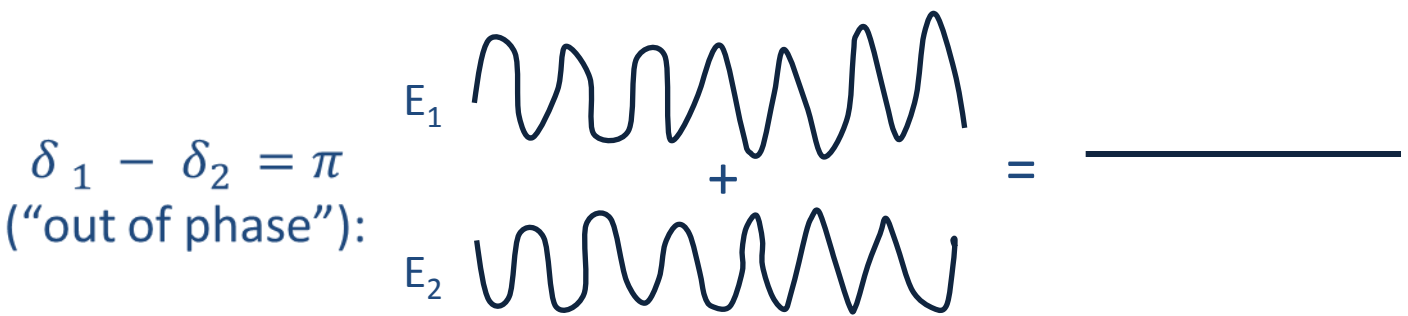
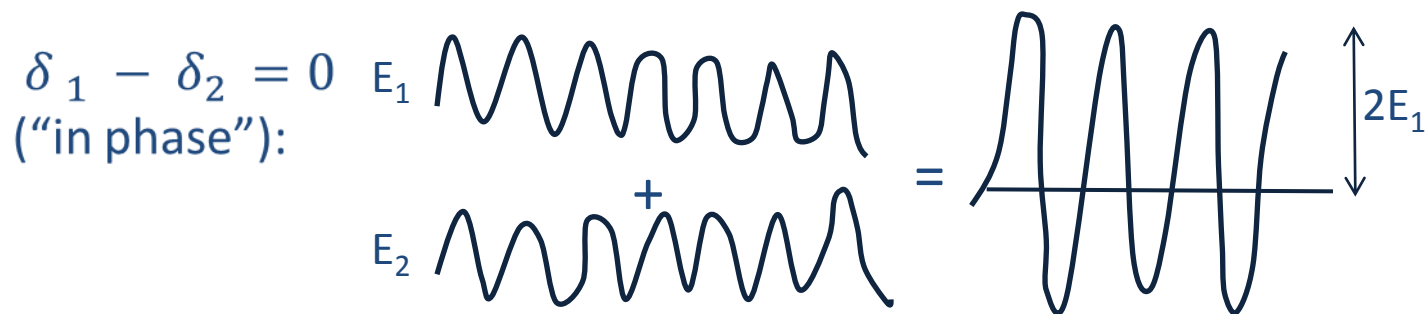
frequency
amplitude
phase

The energy of the wave is proportional to the square of the amplitude:

$$\text{energy} \propto E_0^2$$

# INTERFERENCE

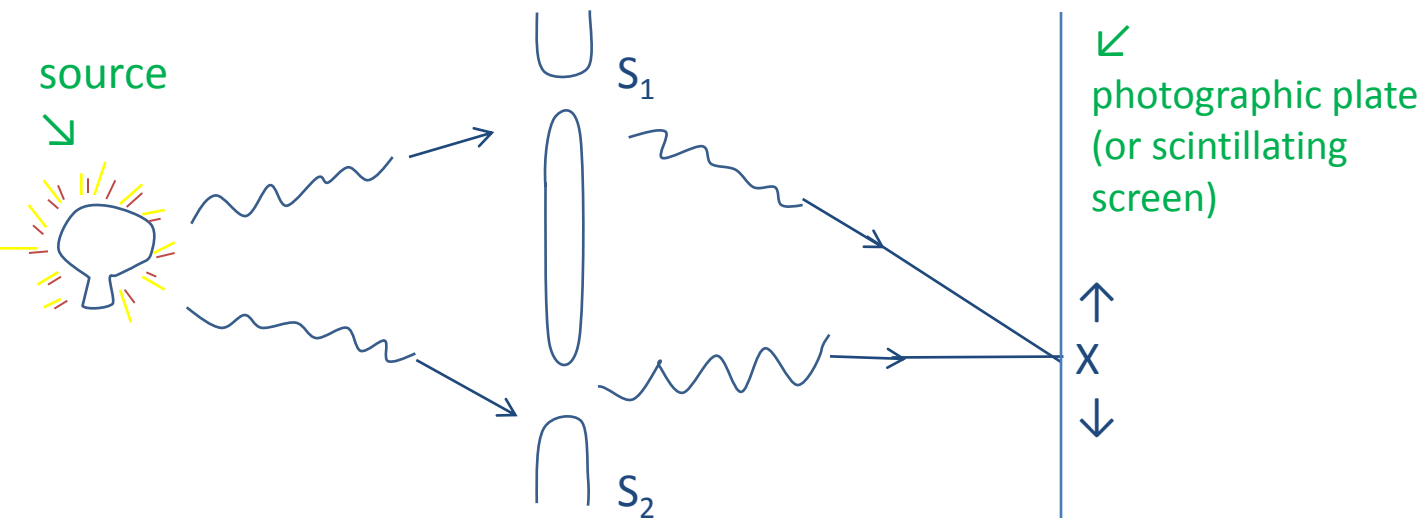
Suppose we try to combine (“superpose”) two light waves. If we sit at a particular point, then at any given time the electric field will be the sum of the fields on the two waves. So it depends on their relative phase: e.g.



If  $|E_1| = |E_2|$ , then energy of 1<sup>st</sup> wave = energy of second wave  $\propto E_1^2$ . But in first case, energy of combined wave  $\propto (2E_1)^2 \propto 4E_1^2$ , and in second case it is 0!

## INTERFERENCE, cont.

Famous example of interference:  
Young's slits:

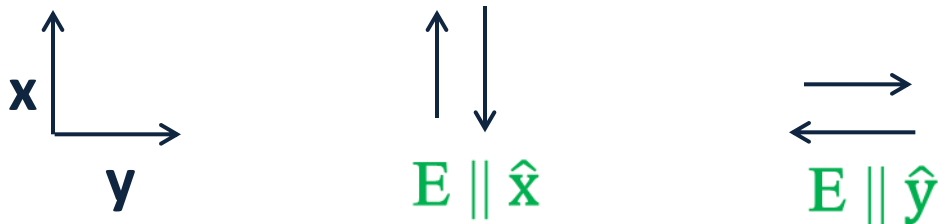


Amplitude  $E$  of total field at  $X$  depends on difference of phase ( $\delta_1 - \delta_2$ ), which in turn depends on difference in path length traversed  $\Rightarrow E$ , hence energy deposited, depends on position of  $X$  on plate  $\Rightarrow$  pattern of light ( $E = \text{max}$ , high-energy) and dark ( $E = 0$ , no energy) bands on plate.

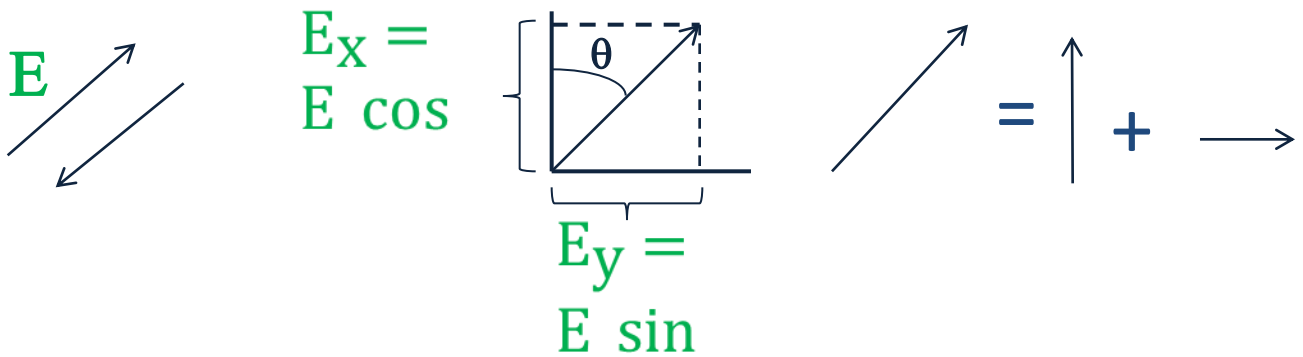
Note: If either slit  $S_1$  or  $S_2$  is closed, electric field  $E$ , and hence distribution of energy on plate is (nearly) uniform (independent of position  $X$ )

## POLARIZATION

The electric field also has a direction associated with it: e.g. for the wave propagating into the screen, we can have



But we can also have e.g.



We can express  $E$  as a “superposition” of its two components (“polarizations”)

$$E = \hat{x} E_x + \hat{y} E_y \equiv E (\hat{x} \cos \theta + \hat{y} \sin \theta)$$

Actually,  $E_x$  and  $E_y$  can have different phases: e.g.

$$E_x = E_0 \cos \omega t, E_y = E_0 \cos(\omega t + \pi/2) = E_0 \sin \omega t$$

Then, if we sit at a particular point, the electric field rotates as a function of time (“circular polarization”). This kind of wave carries angular momentum.

## POLARIZERS

A polarizer is a crystal which (e.g.) transmits a light wave completely when its polarization is vertical ( $||\hat{x}$ ) and reflects it completely when the polarization is horizontal ( $||\hat{y}$ )



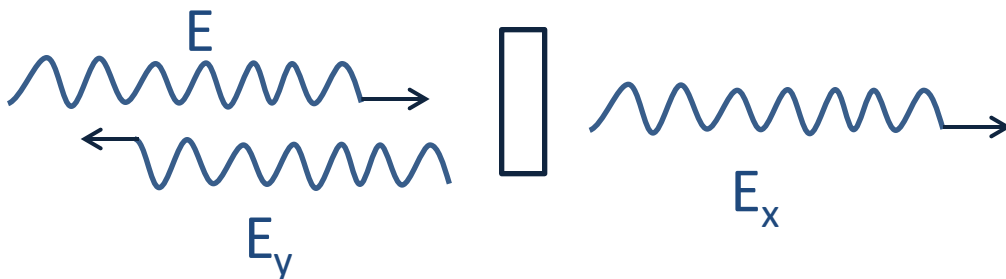
But what if the polarization is at an angle  $\theta$  in the  $xy$ -plane?



Answer: resolve into its components!

$$\mathbf{E} = \hat{x} E_x + \hat{y} E_y \equiv E_0 (\hat{x} \cos \theta + \hat{y} \sin \theta)$$

then the  $x$ -component is completely transmitted and the  $y$ -component completely reflected:



Since energy  $\propto E^2$ ,

$$\text{fraction of energy transmitted} = E_x^2 / E^2 = \cos^2 \theta$$

("Malus's law")

## SUPERPOSITIONS ARE NOT “MIXTURES”!

In the Young's slits experiment, the effect of light coming through both slits is **not** simply the sum of the effects of the two beams coming through the two slits individually.

Another example: transmission through a polarizer (↑ transmitted, →reflected):

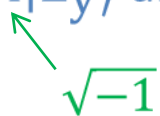
for a beam with polarization  45° 50% of energy is reflected  
( $\cos^2 45^\circ = 1/2$ )

for a beam with polarization  45° 50% is also reflected  
( $\cos^2 45^\circ = 1/2$ )

But, a superposition of  and  is just  =  for which

**no energy is reflected at all!**

So, in dealing with waves, must first **add amplitudes**, and only then calculate energies etc.

Notation (for future reference):  $E_x \equiv$  amplitude of electric field for light wave with vertical ( $\hat{x}$  -) polarization,  $E_y \equiv$  amplitude for wave with horizontal ( $\hat{y}$  -) polarization. Then it turns out that amplitude for RH circularly polarized light wave is  $E_x + i|E_y\rangle$  and state of LH circularly polarized one is  $E_x - i|E_y\rangle$    $\sqrt{-1}$

## QUANTUM MECHANICS: PHOTONS

In quantum mechanics, the energy of a light wave of frequency  $\omega$  comes in **irreducible packets** (p “photons”) of energy  $\hbar\omega$ .

$$h/2\pi, h \equiv \text{Planck's constant}$$

The “amplitude” of a photon is proportional to the amplitude of the electric field of the corresponding classical light wave, so e.g. if

$|x\rangle$  represents amplitude for a photon polarized in x-direction

$|y\rangle$  represents amplitude for a photon polarized in y-direction

Then

$\cos\theta|x\rangle + \sin\theta|y\rangle \equiv |\mathbf{a}\rangle$  represents amplitude for photon polarized in direction  $\mathbf{a}$  in xy-plane, making  $\angle\theta$  with x-axis, and

$|x+iy\rangle$  represents amplitude for RH circularly polarized photon

$|x-iy\rangle$  represents amplitude for LH circularly polarized photon

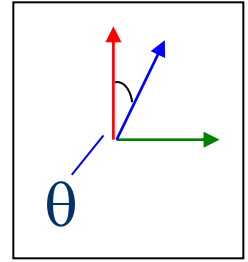
Note: RH (LH) circularly polarized photons carry angular momentum  $\hbar(-\hbar)$

Two-photon processes: Suppose an atom emits two photons with definite polarizations, say  $|\mathbf{a}\rangle$  for 1 and  $|\mathbf{b}\rangle$  for 2. Then the combined state is described by an **amplitude** which is the product of the amplitude for the two photons separately:

$$\text{amplitude} = |\mathbf{a}\rangle_1 |\mathbf{b}\rangle_2$$



Single photon incident on  
birefringent crystal (“polarizer”):

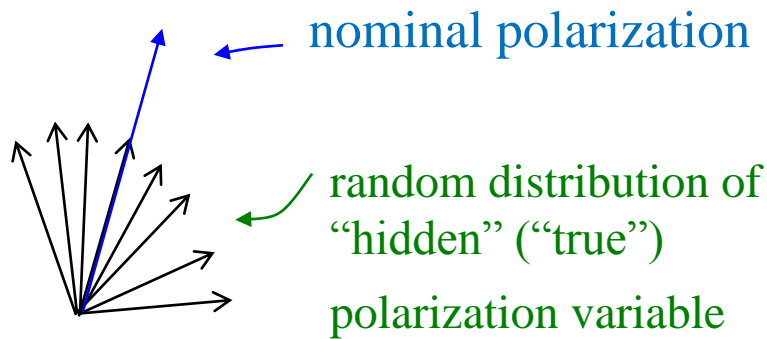


**Probability** of transmission =  $\cos^2\theta$

(quantum version of Malus' law)

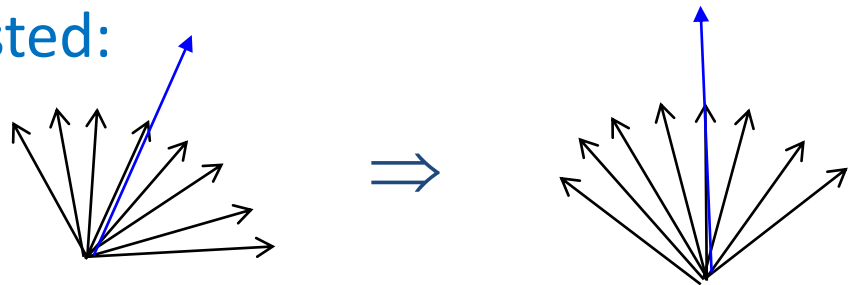
Digression: Can a classical probabilistic  
theory explain this?

**YES!**



“true” polarization

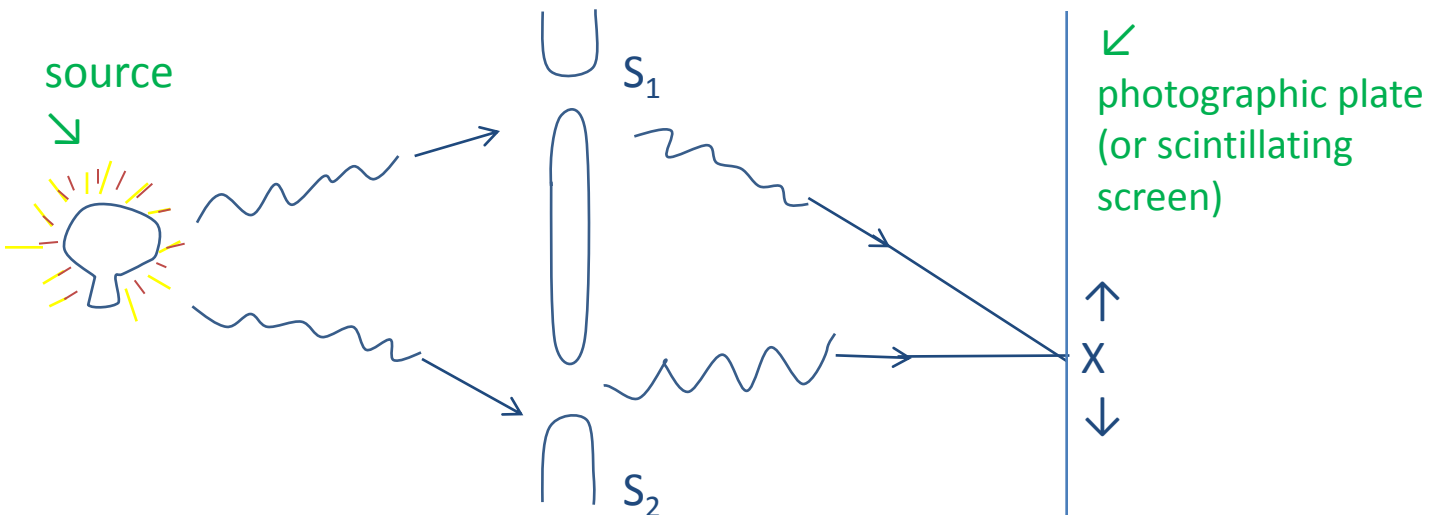
If closer to , transmitted: if closer to , reflected. If transmitted, distribution of “hidden” variable is adjusted:



With suitable choice of random distribution,

can reproduce  $P_T(\theta) = \cos^2 \theta$      $f(\chi) = \cos 2(\theta - \chi)$

# PHOTONS IN THE YOUNG'S SLITS EXPERIMENT



If we turn down the strength of the source until photons come through one at a time, we can see them arriving individually on the screen. Initially, arrival seems random, but eventually a pattern builds up, with

**probability of arrival at  $X \propto$  brightness of classical interference pattern**

Since classically the brightness  $\propto$  energy  $\propto E^2$  this suggests that in quantum mechanics

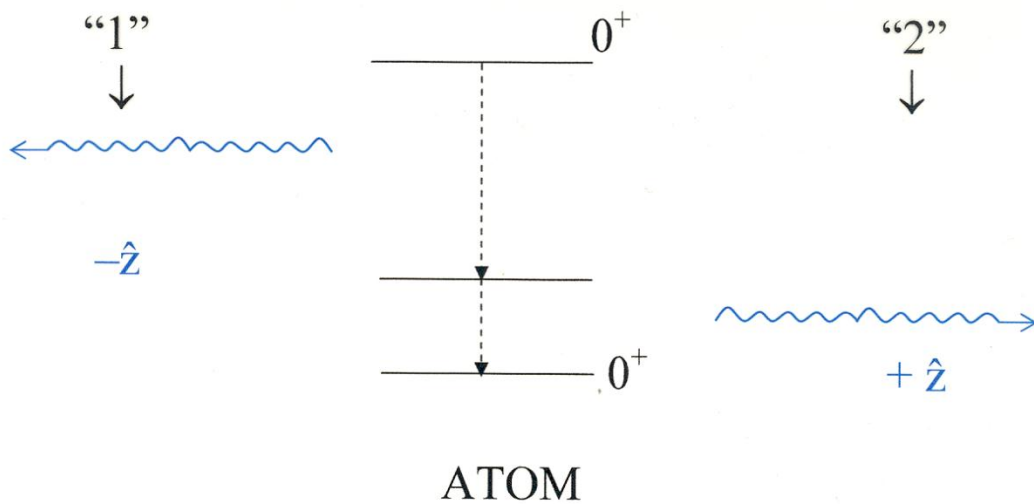
$$\text{PROBABILITY} \propto (\text{AMPLITUDE})^2$$

NOTE: If either slit closed, probability is uniform (ind. of  $X$ ). Thus in QM, cannot add probabilities of two mutually exclusive events (paths), but

**must add amplitudes not probabilities!**

Couldn't we arrange to detect which slit a particular photon came through? Yes (in principle), but then we destroy the interference pattern! (i.e. prob. of arrival is ind. of  $X$ )

## 2-PHOTON STATES FROM CASCADE DECAY OF ATOM



What is polarization state of photons?

(note:  $|x\rangle \pm i|y\rangle$  corr. photon angular momentum  $\pm \hbar$ )

General principle of QM; if process can happen either of two ways, **and we don't (can't) know which**, must add amplitudes!

Here, we know that total angular momentum of 2 photons is zero, but we don't know whether photon 1 carried off  $+\hbar$  and 2  $-\hbar$  (intermediate atomic state has  $m = 1$ ) or vice versa (int. state  $m = -1$ ). Hence must write

$$|\psi\rangle = |x + iy\rangle_1 |x - iy\rangle_2 + (\text{phase factor}).$$

$$|x - iy\rangle_1 |x + iy\rangle_2$$

**crucial!**

Actually (parity  $\Rightarrow$ ) phase factor = + 1, so

$$|\psi\rangle = |x\rangle_1 |x\rangle_2 + |y\rangle_1 |y\rangle_2 \quad \left( x \frac{1}{\sqrt{2}} \right)$$

## POLARIZATION STATE OF 2 PHOTONS EMITTED BACK TO BACK IN ATOMIC $0^+ \rightarrow 1^- \rightarrow 0^+$ TRANSITION (recap):

$$\Psi_{2\gamma} = \frac{1}{\sqrt{2}} (|x\rangle_1 |x\rangle_2 + |y\rangle_1 |y\rangle_2)$$

So, if photon 1 is measured to have polarization  $x(y)$  so inevitably will photon 2!

But, state is rotationally invariant:



$$\Psi_{2\gamma} = \frac{1}{\sqrt{2}} (|x'\rangle_1 |x'\rangle_2 + |y'\rangle_1 |y'\rangle_2)$$

so if 1 measured to have polarization  $x'$  ( $y'$ ) so does 2!

(NOT true for “mixture” of  $|x\rangle_1 |x\rangle_2$  and  $|y\rangle_1 |y\rangle_2$ )

Now:

What if photon 1 is incident on polarizer with “transmission” axis  $\hat{a}$ , and photon 2 on one with a differently oriented transmission axis  $\hat{b}$ ?

Since choice of axes for  $\Psi_{2\gamma}$  arbitrary, choose  $\hat{x} = \hat{a}$ .

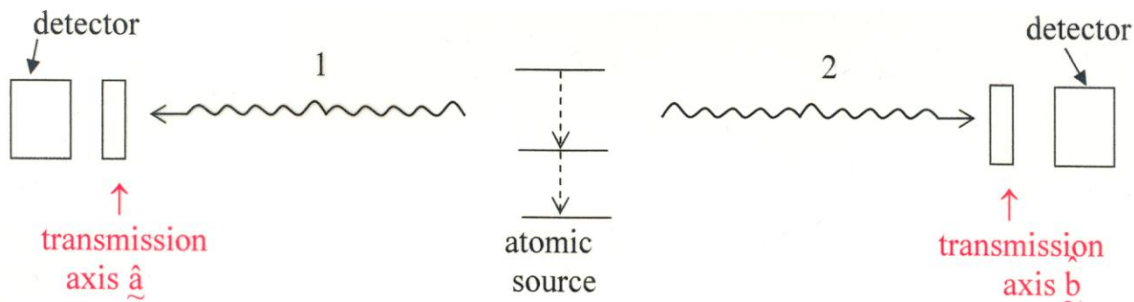
Then:

Prob. of transmission of 1 =  $\frac{1}{2}$ .

But if 1 transmitted, then polarization of 2 is  $\hat{a}$ , so probability of transmission of polarizer set in direction  $\hat{b} = \cos^2 \theta_{ab}$

(Malus’s law)

$$P(\text{both transmitted}) = \frac{1}{2} \cos^2 \theta_{ab}$$



$$\text{Prob. (both detected)} = \frac{1}{2} \cos^2 \theta_{ab} = \frac{1}{4} (1 + \cos 2\theta_{ab})$$

$$\text{Prob. (neither detected)} = \frac{1}{2} \cos^2 \theta_{ab}$$

$$\left. \begin{array}{l} \text{Prob. (1 detected, 2 not)} \\ \text{Prob. (2 detected, 1 not)} \end{array} \right\} = \frac{1}{2} \sin^2 \theta_{ab}$$

“Isotropic mixture”  
 $\Rightarrow \frac{1}{4} (1 + \frac{1}{2} \cos 2\theta_{ab})$

THESE ARE THE PREDICTIONS OF **STANDARD QUANTUM MECHANICS**. CAN THEY BE EXPLAINED BY A CLASSICAL PROBABILISTIC THEORY?

Df: If for a given pair, with polarizer 1 set at  $\hat{a}$ , photon 1 is detected,

df.  $A \equiv +1$ : if rejected,  $A \equiv -1$ . Similarly with polarizer 2 set at  $\tilde{b}$ , if

photon 2 detected, df.  $B \equiv +1$ : if rejected, then  $B \equiv -1$ . Then above is equivalent to the statement that for the average over the ensemble of pairs,

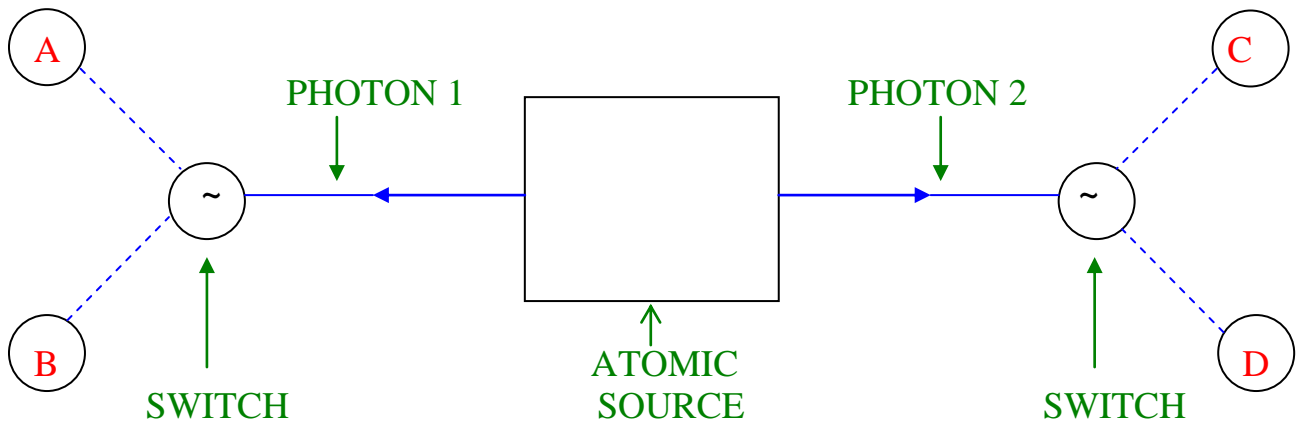
$$\langle AB \rangle = \cos^2 \theta_{ab} - \sin^2 \theta_{ab} = \cos 2\theta_{ab}$$



[“Mixture”:  $\frac{1}{2} \cos 2\theta_{ab}$ ]

as special cases, for  $\theta_{ab} = 0$   $\langle AB \rangle = +1$ , and for  $\theta_{ab} = \pi/2$ .  $\langle AB \rangle = -1$ .

(EPR). These two special cases can be accounted for by a classical probabilistic model. But...

# EXPERIMENTS ON CORRELATED PHOTONS



(A)  $\equiv$   , etc.)  
 ↗ transm. axis =  $\tilde{a}$

DEFINITION: If photon 1 is switched into counter “A”, then:

If counter “A” clicks,  $A = +1$  (DF.)

If counter “A” does not click,  $A = -1$  (DF.)

NOTE:

If photon 1 switched into counter “B”, then A is NOT DEFINED.

Experiment can measure

$\langle AC \rangle_{\text{exp}}$  on one set of pairs ( $1 \rightarrow$  “A”,  $2 \rightarrow$  “C”)

$\langle AD \rangle_{\text{exp}}$  on another set of pairs ( $1 \rightarrow$  “A”,  $2 \rightarrow$  “D”)

etc.

Of special interest is

$$K_{\text{exp}} \equiv \langle AC \rangle_{\text{exp}} + \langle AD \rangle_{\text{exp}} + \langle BC \rangle_{\text{exp}} - \langle BD \rangle_{\text{exp}}$$

for which Q.M. makes clear predictions.

## POSTULATES OF “OBJECTIVE LOCAL” THEORY:

- (1) Local causality
- (2) Induction
- (3) Microscopic realism OR macroscopic  
“counter-factual definiteness”

## BELL’S THEOREM

1. (3)  $\rightarrow$  For each photon 1, EITHER  $A = +1$  OR  $A = -1$ , independently of whether or not A is actually measured.
2. (1)  $\rightarrow$  Value of A for any particular photon 1 unaffected by whether C or D measured on corresponding photon 2. : etc.
3.  $\therefore$  For each pair, quantities AC, AD, BC, BD exist, with A, B, C, D,  $= \pm 1$  and A the same in (AC, AD) (etc.)
4. Simple algebra then  $\rightarrow$  for each pair,  $AC + AD + BC + -BD \leq 2$
5. Hence for a single ensemble,  

$$\langle AC \rangle_{\text{ens}} + \langle AD \rangle_{\text{ens}} + \langle BC \rangle_{\text{ens}} - \langle BD \rangle_{\text{ens}} \leq 2$$
6. (2)  $\rightarrow \langle AC \rangle_{\text{exp}} = \langle AC \rangle_{\text{ens}}$ , hence the measurable quantity  

$$K_{\text{exp}} = \langle AC \rangle_{\text{exp}} + \langle BC \rangle_{\text{exp}} + \langle BC \rangle_{\text{exp}} - \langle BD \rangle_{\text{exp}}$$
satisfies

$$K_{\text{exp}} \leq 2, \text{ Obj. Local Theory}$$



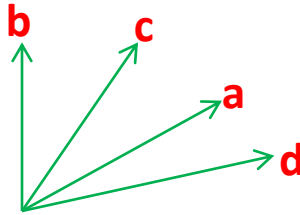
OBJECTIVE LOCAL THEORY:  $K_{\text{exp}} \leq 2$ .

QM: If polarizer settings are **a**, **b**, **c**, **d**

Then e.g. for a  $0^+$  transition predict

$$\langle AC \rangle = \cos(2\theta_{\underline{a} \cdot \underline{c}}), \text{ etc.}$$

$\Rightarrow$  for



QM predicts (ideal case)

$$\underline{K_{\text{exp}} = 2\sqrt{2}}$$

$\Rightarrow$  **Exptl. Predictions** of QM incompatible with those of an theory embodying

- Local causality
- Induction
- Macroscopic counter-factual definiteness

1. “It is a fact that either A would have clicked or A would not have clicked”
2. “Either it is a fact that A would have clicked, or it is a fact that A would not have clicked”