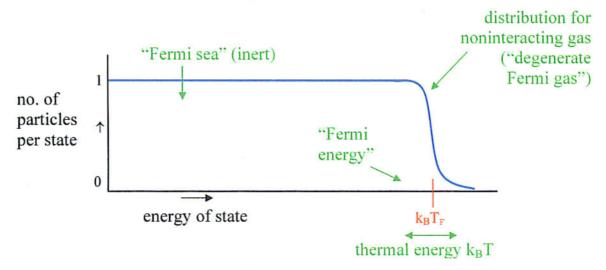
## SUPERFLUID <sup>3</sup>HE: SOME PRE-HISTORY

<sup>4</sup>He: below 2K, superfluid



Particles of spin  $\frac{1}{2}$   $\Rightarrow$  Fermi statistics



Landau (1957): interactions don't change picture qualitatively (in "normal" phase) ("degenerate Fermi liquid")

$$T_F \sim 10^4 - 10^5 \, \mathrm{K}$$
  $T_F \sim 5 \, \mathrm{K}$  at  $T \lesssim 20 \, \mathrm{K}$ . at  $T \lesssim 10^{-3} \, \mathrm{K}$ . superconductivity (in some metals)

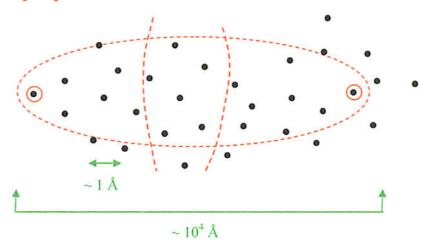
#### MORE PREHISTORY

Theory of superconductivity:

(a) phenomenological (V. L. Ginzburg, A. A. Abrikosov, et al., 1950-1955):

macroscopic wave function

(b) microscopic (Bardeen et al., 1957): temp.,  $\lesssim 20 \mathrm{K}$  electrons in energy shell of width  $\sim k_B T_C$  around Fermi energy form Cooper pairs



critical

Crucial feature of BCS theory: ALL COOPER PAIRS MUST BEHAVE IN EXACTLY THE SAME WAY!

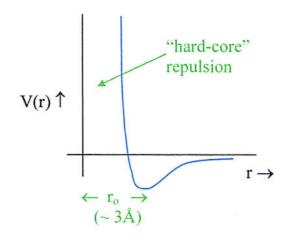
(GL "macroscopic wave function" is just the **common** center-of-mass wave function of all the pairs)

In BCS theory, "internal" wave function of pairs trivial: (" <sup>1</sup>S<sub>o</sub>")

$$\psi \text{ } (\underline{r}_1\underline{r}_2 : \sigma_1\sigma_2) \sim \frac{1}{\sqrt{2}} \text{ } (\uparrow_1 \downarrow_2 - \downarrow_1 \uparrow_2) \text{ } f \text{ } (|\underline{r}_1 - \underline{r}_2|)$$
 spherically symmetric spin singlet 
$$(\ell = 0)$$

NO INTERNAL ("ORIENTATIONAL") DEGREES OF FREEDOM

# EARLY THEORETICAL WORK ON POSSIBLE COOPER PAIRING IN LIQUID <sup>3</sup>HE



$$r \sim r_0, \ p \sim p_F \ (\equiv \sqrt{2mk_BT_F} \ )$$
  
 $\Rightarrow$  relative angular momentum  
 $\ell \equiv (p_F r_0 / \ \hbar) \neq 0$   
(prob. 1 or 2)

Pauli principle: 
$$\begin{cases} \ell = 0, 2, 4... & S = 0 \text{ (singlet)} \\ \ell = 1, 3, 5... & S = 1 \text{ (triplet)} \end{cases}$$

in general,  $\ell \neq 0 \Rightarrow$  relative (internal) wave function of pair has orientational degree(s) of freedom! "equal spin pairing"

Anderson & Morel (1961): explore in detail case  $\ell = 2$ , and a special case of  $\ell = 1$ : only  $\uparrow \uparrow$  and  $\downarrow \downarrow$  pairs form, and have the same orbital ang. momentum in direction  $\hat{\ell}$  ("ABM" state) Physical properties anisotropic.

Vdovin
Balian & Werthamer (1963): in  $\ell = 1$  case all spin components

( $\uparrow \uparrow, \downarrow \downarrow, \frac{1}{\sqrt{2}} \uparrow \downarrow + \downarrow \uparrow$ )) can form: in fact for any given pair,  $L = -S \Rightarrow J = 0$ .

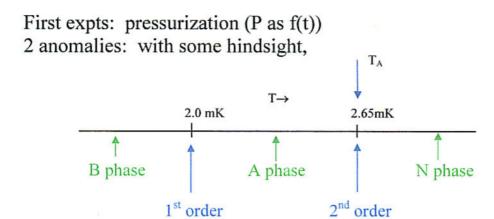
("BW" state). All physical properties isotropic. More stable than any ESP state.

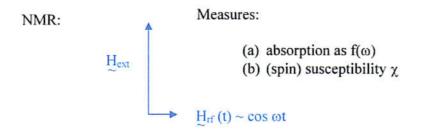
Theoretical expectation c. 1964:

Liquid <sup>3</sup>He may form Cooper pairs, either with  $\ell$  = even (spin singlet) or with  $\ell$  = odd (BW state). In either case,  $\chi$  reduced and all magnetic properties isotropic.  $T_c$  difficult to predict.

# THE EXPERIMENTS OF 1971-72 (D. D. Osheroff, R. C. Richardson, D. M. Lee...(Nobel prize 1996)):

Mixture of liquid and solid  ${}^{3}$ He, T < 3 mK. (so only temperature varied).



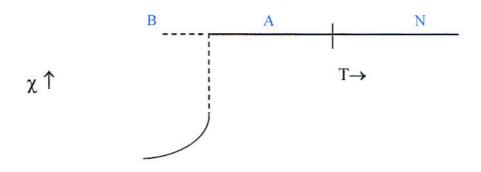


#### In N state:

χ independent of temperature, value as expected for degenerate Fermi liquid

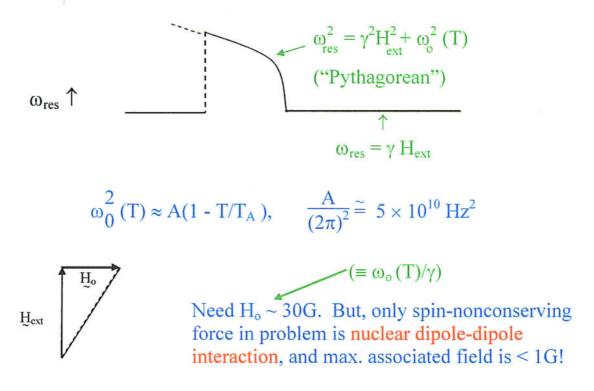
Abs<sup>n</sup> shows v. sharp peak at free-atom Larmor frequency:  $\omega_{res} = \gamma H_{ext}$  gyromagnetic ratio of free <sup>3</sup>He atom, ~ 3000 Hz/G

#### NMR in the new phases:



Not necessarily mysterious: e.g. A phase could be an ESP state (only  $\uparrow\uparrow,\downarrow\downarrow$  pairs  $\Rightarrow$  no reduction in  $\chi$ ), B could be singlet or BW (some  $\uparrow\downarrow$  pairs, so  $\chi$  reduced) [but: why is ESP ever stable?] -

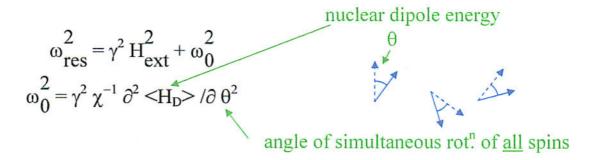
But: what about the resonance frequency?



IS THIS THE FIRST INDICATION OF A RADICAL BREAKDOWN OF QUANTUM MECHANICS?

#### WHAT CAN BE INFERRED FROM SUM RULES?

IF a single sharp resonance is observed (as in expt.) then:



But 
$$\partial^2 \langle H_D \rangle / \partial \theta^2 \sim \langle H_D \rangle$$
:  
So, exptl. value of  $\omega_0^2$  (T)  $\Rightarrow$ 

$$\langle H_D \rangle$$
 (T) ~ K(1 - T/T<sub>A</sub>), K ~ 10<sup>-3</sup> ergs/cm<sup>3</sup>

HOW CAN THIS BE?

$$\begin{cases} \uparrow \text{ ("bad")} & \uparrow \\ \rightarrow \text{ ("good")} & \rightarrow \end{cases}$$
 
$$\Delta E \lesssim \frac{\mu_o \mu_n^2}{\frac{3}{r_o^2}} \sim 10^{-7} \text{ K } \ll k_B T$$

So, prima facie, preference for "good" orientation over "bad" is at most

$$\sim \Delta E/k_BT \sim 10^{-4}$$
 [actually,  $\sim \Delta E/k_BT_F \sim 10^{-7}$ ]

⇒ expectation value of dipole energy much too small!

#### SPONTANEOUSLY BROKEN SPIN-ORBIT SYMMETRY

#### Ferromagnetic analogy:

#### **FERROMAGNET**

$$\stackrel{\wedge}{H}=\stackrel{\wedge}{H_o}+\stackrel{\wedge}{H_z}$$

1

invariant under simult. rotation of all spins

extl. field 
$$\hat{H}_z = -\mu_B \mathcal{H} \sum_i S_{zi}$$
 breaks spin-rot. symmetry

Paramagnetic phase (T > T<sub>c</sub>): spins behave independently, kT competes with  $\mu_B \mathcal{H} \Rightarrow$ polarization  $\sim \mu_B \mathcal{H}/kT \ll 1 \Rightarrow$  $<H_z> \sim N(\mu_B \mathcal{H})^2/kT$ 

Ferromagnetic phase (T < Tc):  $\dot{H}_o$  forces all spins to lie parallel  $\Rightarrow k_B T$  competes with  $N\mu_B \mathcal{H}$  $\Rightarrow \langle S_z \rangle \sim 1 \Rightarrow \langle H_z \rangle \sim N\mu_B \mathcal{H}$ 

## LIQUID 3HE

$$\overset{\,\,{}_{}}{H}=\overset{\,\,{}_{}}{\overset{\,\,{}_{}}{H}}{}_{o}+\overset{\,\,{}_{}}{H}{}_{D}$$

invariant under <u>relative</u> rotation of spin + orbital coordinate systems

$$= \mu_o \mu_n^2 / r_o^3$$

$$\hat{H}_D = g_D \sum_{ij} \left( \frac{g_i \cdot g_j - 3g_i \cdot \hat{\underline{r}}_{ij} g_j \cdot \hat{\underline{r}}_{ij}}{(r_{ij}^3 / r_o^3)} \right)$$

breaks relative spin-orbit rot." symmetry

Normal phase  $(T > T_A)$ :

pairs of spins behave

independently  $\Rightarrow$ polarization  $\sim g_D/kT \ll 1 \Rightarrow$   $< H_D > \sim N g_D^2/kT$ 

Ordered phase  $(T < T_A)$ :  $H_o$  forces all pairs to

behave similarly  $\Rightarrow$  kT competes with  $Ng_D$   $\Rightarrow <H_D> \sim Ng_D$   $\sim 10^{-3} \text{ ergs/cm}^{3}$ !

#### SBSOS: ORDERING MAY BE SUBTLE

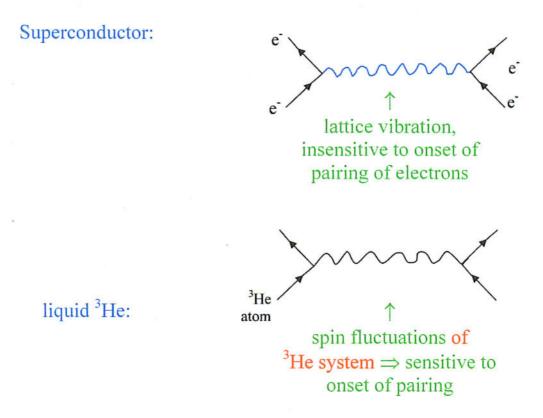
# LIQUID 3HE **FERROMAGNET NORMAL PHASE** ORDERED **PHASE** ( $\nearrow$ = total spin of pair = relative orbital ang. momentum) $\langle S \rangle = \langle L \rangle = 0$ $\langle \mathbf{S} \rangle \neq 0$ but $\langle L \times S \rangle \neq 0!$

#### RESOLUTION OF THE PARADOX OF TWO NEW PHASES.

(Anderson & Brinkman, Phys. Rev. Letters 30, 1108 (1973))

In BCS (weak-coupling) theory for  $\ell=1$ , BW phase is always stable, independently of pressure and temperature.

Crucial difference between Cooper pairing in superconductors and <sup>3</sup>He



⇒ "feedback" effects: Over most of the phase diagram, BW state stable as in BCS theory. But at high temperature and pressure, feedback effects uniquely favor ABM phase.

major qualitative leap beyond BCS!

#### MICROSCOPIC SPIN DYNAMICS (SCHEMATIC)

Basic variables:

Total spin S (a)

(b)

$$\label{eq:continuous} \text{Orientation } \underbrace{\emptyset} \text{ of spin of Cooper pairs} \qquad \left\{ \begin{aligned} [S_i,\,\theta_j] &= i\delta_{ij} \end{aligned} \right.$$

$$\hat{\mathbf{H}} = \hat{\mathbf{H}}_{o}(\mathbf{S}) + \hat{\mathbf{H}}_{D}(\mathbf{\theta})$$

$$\uparrow$$

hydrodynamic (Born-Oppenheimer) approximation

Semiclassical equations of motion:

dipole torque

$$\frac{d \underbrace{\theta}}{dt} = \frac{\partial < \dot{H}_o>}{\partial \underbrace{S}} = \mathcal{H}_{ext} - \chi^{-1} \underbrace{S}, \qquad \frac{d \underbrace{S}}{dt} = \underbrace{S} \times \mathcal{H}_{ext} - \frac{\partial < \dot{H}_o>}{\partial \underbrace{\theta}}$$

⇒ linear NMR behavior completely determined by eigenvalues of quantity

$$\Omega_{ij}^2 \equiv \partial^2 \langle H_D \rangle / \partial \theta_i \partial \theta_j$$
 So, can Tingerprint A and B phases by

so, can "fingerprint" NMR!

ABM: single resonance line

axial: split resonance

BW: original BW state is  $\underline{L} = -\underline{S}$ , i.e. J = 0. But dipole torque rotates S relative to L by  $\angle \cos^{-1}(-1/4) = 104^{\circ}$  around axis  $\hat{\omega}$  whose "best" choice is  $\mathcal{H}_{\rm ext}$ .

Result: no shift in transverse resonance, but finite-frequency longitudinal resonance!  $\mathcal{H}_{\text{ext}}$ (also in ABM phase)

#### CONCLUSION (by summer of 1973):

Both a priori stability considerations and NMR experimental data are consistent with hypothesis that both new phases are Cooper-paired ("superfluid") phases. Specifically,

```
A phase = ABM
B phase = BW
```

### What is superfluid <sup>3</sup>He good for?

- (a) most sophisticated physical system of which we can claim detailed quantitative understanding. E.g. textures, orientational dynamics, topological singularities...
- (b) analogies with systems in particle physics, cosmology... (G. E. Volovik)
- (c) studies of (some aspects of) turbulence
- (d) Amplification of ultra-weak effects (cf NMR): Example: P- (but not T-) violating effects of neutral current part of weak interaction:

For single elementary particle, any EDM d must be of form

$$\underline{d} = \text{const. } \underline{J} \leftarrow \text{violates T as well as P.}$$

violates P but not T.

Effect is tiny for single pair, but since all pairs have same value of  $\mathbb{L} \times \mathbb{S}$ , is multiplied by factor of  $\sim 10^{23} \Rightarrow$ 

macroscopic P-violating effect?