

TOPOLOGICAL QUANTUM COMPUTING
IN $(p+ip)$ FERMI SUPERFLUIDS
AND STRONTIUM RUTHENATE:
PROSPECTS AND PROBLEMS

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BRIEF REVIEW OF “ESTABLISHED WISDOM”

1. p + ip Fermi superfluids

order parameter

For a general spin-1/2 Fermi superfluid, OP df. by

$$F_{\alpha\beta}(\underline{r}, \underline{r}') \equiv \langle \psi_{\alpha}^{\dagger}(\underline{r}) \psi_{\beta}^{\dagger}(\underline{r}') \rangle \leftarrow \text{"anomalous average"}$$

p + ip:

$$F_{\alpha\beta}(\underline{r}, \underline{r}') = \delta_{\alpha\beta} F_{\alpha}(\underline{r}, \underline{r}') \quad (\text{ESP})$$

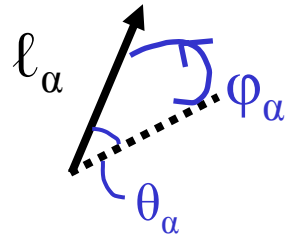
↑ “equal spin pairing”

$$F_{\alpha\beta}(\underline{r}, \underline{r}') \equiv F_{\alpha}(\underline{R}, \underline{\rho}) \cong F_{\alpha}(\underline{R}) f_{\alpha}(\underline{\rho})$$

COM → ↑ ← rel. coord.

$$f_{\alpha}(\underline{\rho}) \equiv \sin \theta_{\alpha} \exp i\varphi_{\alpha}$$

↑
breaks TRI



thus if ℓ_{α} taken as z-axis, F. T. is

$$F_{\alpha}(\underline{p}) = p_x + ip_y \leftarrow \text{"p + ip"}$$

Standard ansatz for MBWF (COM's of \uparrow, \downarrow at rest):

$$\Psi = \Psi_{\uparrow} \Psi_{\downarrow}$$

$$\Psi_{\alpha} \equiv \left(\sum_{\mathbf{k}} c_{\mathbf{k}}^{\alpha} a_{\mathbf{k}_{\alpha}}^{+} a_{-\mathbf{k}_{\alpha}}^{+} \right)^{N/2} | \text{vac} \rangle$$

where, if z-axis chosen along ℓ_{α} .

$$c_{\mathbf{k}}^{\alpha} = f(|\mathbf{k}|) \cdot (k_x + ik_y) \sim \exp i\varphi_{\mathbf{k}}$$

Note dependence on $\varphi_{\mathbf{k}}$ extends to whole Fermi sea

Some properties of “standard” ansatz:

1. Ang. Momentum along $\ell_{\alpha} \cong N_{\alpha} \hbar / 2 (-0(\Delta / E_F))$
even in limit $\Delta \rightarrow 0$ or $T \rightarrow T_c$.
2. For 2D (“planar”) case ($\ell_{\alpha} \perp$ plane), put $z = x + iy$, then for all $|z_i - z_j| \gg \xi$, coord-space MBWF is of form

↑
pair radius

$$\Psi_N(z_1 z_2 \dots z_N) = \text{const. } \underbrace{\mathcal{A} \left\{ \frac{1}{z_1 - z_2} \cdot \frac{1}{z_3 - z_4} \cdot \frac{1}{z_5 - z_6} \dots \right\}}_{\text{“Pfaffian”}}$$

Cf MR ansatz for $\nu = 5/2$ QH state

↑
Moore-Read



ESTABLISHED WISDOM (cont.)

Examples of (p + ip) Fermi superfluids:

1. $^3\text{He-A}$

Evidence:

- (a) ESP: χ unchanged, NMR
- (b) p-wave: sp. ht., ultrasound attenuation
- (c) broken TRI: (literal) ferromagnetism (?)
- (d) specific (p + ip) state (3D): NMR

2. Sr_2RuO_4 (“SRO”)

(structure similar to cuprates, $T_c \sim 1.5$ K)

Evidence:

- (a) ESP: χ unchanged in sup^s phase
- (b) p-wave-like (i.e. odd-parity): Josephson (PSU)
- (c) broken TRI: muon spin resonance, Josephson (UIUC), Kerr rotation (but \uparrow : also in N state of cuprates)



ESTABLISHED WISDOM (cont.)

3. Half-quantum vortices (“HQV”)

Should occur in any ESP Fermi superfluid, provided coupling between $\uparrow\uparrow$ and $\downarrow\downarrow$ sufficiently weak.

e.g. (neutral case):

vortex in $\uparrow\uparrow$ components, nothing in $\downarrow\downarrow$ component, i.e.

$\Delta_{\uparrow\uparrow} \propto \exp i\Phi$, $\Delta_{\downarrow\downarrow} \propto \text{const.}$ (“half-quantum” vortex)

Note, however, that quantization condition for $\uparrow\uparrow$ pair velocity is still

$$K \equiv \int_0^{2\pi} v_{s\uparrow} \cdot dl = h / 2m$$

Can tolerate Majorana fermions.

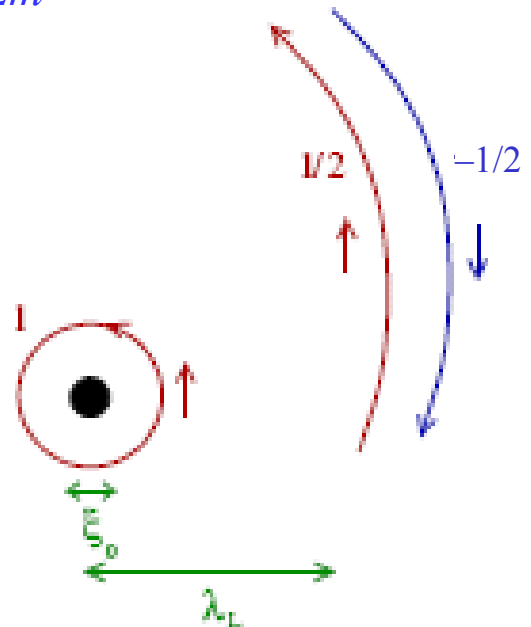
What about charged case?

At $r \ll \lambda_L$, current only in $\uparrow\uparrow$ component \rightarrow total $j \neq 0$.

However, for $r \gtrsim \lambda_L$, $\downarrow\downarrow$'s are involved :

$$0 \neq \underline{j}_{\uparrow} = \underline{j}_{\downarrow}$$

$$\underline{j}_{\text{tot}} = 0$$



Total trapped flux = $\Phi_0/2 = nh/4e$.

Hence, vortex with Dirac (non-Majorana) fermion circling another picks up phase $\pi/2$, not π as for BCS)

Note: HQV in charged system carries circulating spin current as $r \rightarrow \infty \Rightarrow$ energetically disadvantaged relative to simple $(h/2e)$ vortex.

ESTABLISHED WISDOM (cont.)

4. Majorana Fermions and TQC

← topological quantum computation

Consider a HQV in a **2D** ($p + ip$) Fermi superfluid. Model

“gap function” $\Delta_{\uparrow}(\underline{r}, \underline{r}') (\equiv V(\underline{r} - \underline{r}') F_{\uparrow}(r r'))$

by

$$\Delta_{\uparrow}(\underline{r}, \underline{r}') \equiv \Delta_{\uparrow}(\underline{R}, \underline{\rho}) = \Delta(\underline{R})(\nabla_x + i\nabla_y)\delta(\underline{\rho})$$

COM $\xrightarrow{\uparrow}$ $\xleftarrow{\uparrow}$ rel. coord

Then (Kopnin-Salomaa, Volovik, Moore-Read) \exists a

single solution of the BdG equations for the particle/hole

\uparrow
Bogoliubov-deGennes

amplitudes $u(\underline{r}) / v(\underline{r})$ s.t.

$$E = 0, u(\underline{r}) = v^*(\underline{r})$$

“Majorana fermion” (M.F.)

These M.F.’s satisfy the braiding and recombination rules of **Ising anyons**, and that can be used (Ivanov) for (partially) topologically protected quantum computation, by braiding the HQV’s appropriately.

[End of Established Wisdom]



SOME QUESTIONS ABOUT THE ESTABLISHED WISDOM

1. Nature of MBWF of (p + ip) Fermi superfluid

Recap: standard ansatz is (for say $\uparrow\uparrow$)

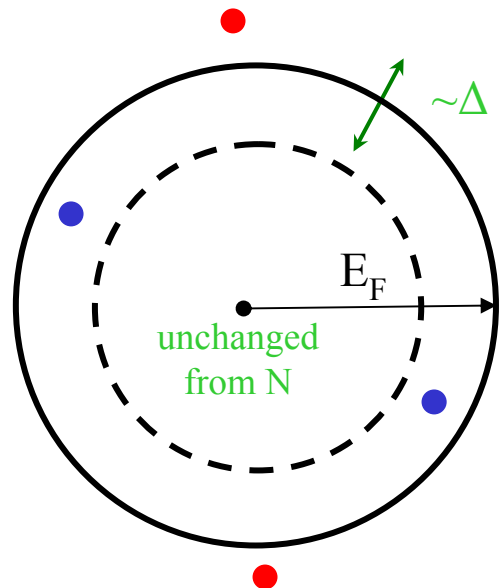
$$\Psi \sim \left(\sum_k c_k a_k^+ a_{-k}^+ \right)^{N/2} | \text{vac} \rangle, c_k \sim \exp i\varphi_k$$

i.e. **all** pairs of states in Fermi sea have anyon momentum \hbar .

Alternative ansatz:

first shot:

$$\Psi(N_p, N_h) \sim \left(\sum_{k > k_F} c_k a_k^+ a_{-k}^+ \right)^{N_p/2} \left(\sum_{k < k_F} d_k a_{-k} a_k \right)^{N_h/2} | \text{vac} \rangle$$



\uparrow : keeps $pp \rightarrow pp$ and $hh \rightarrow hh$, but not (e.g.) $pp \rightarrow hh$.

Remedy:

$$\Psi \sim \sum_{N_p, N_k} Q_{N_p N_k} \Psi(N_p, N_k),$$

Q slowly varying as $f(N_p, N_k)$

degenerate with standard ansatz to $0(N^{-1/2})$, but

$$L \sim (N\hbar / 2) \cdot (\Delta / E_F)^2$$

IS GS OF (p + ip) UNIQUE?



QUESTIONS (cont.)

2. Are the physical systems suitable?

$^3\text{He-A}$

Almost certainly (p + ip), but can it be made “2D”?

e.g. “slab” geometry:



↑: en. gap has modes along $\tilde{\ell}$

⇒ low-lying excitations (v. bad for TQC!)

To eliminate these, need $k_F d \ll \varepsilon_F / k_B T$. But also need $d \gg \xi$. Compatible for $T \ll T_c$.

SRO

Ev. for ESP and violation of TRI fairly strong, but not entirely obvious that latter is due to (p + ip) (e.g. μSR signal could be due to nonunitary spin state?)

In any case, is it sufficiently “2D”?

(a) single layer?

(b) thin macroscopic slab?



QUESTIONS (cont.)

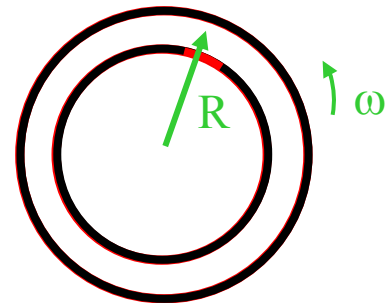
3. Half-quantum vortices

Problems:

(A) Non-observation of HQV's in $^3\text{He-A}$:

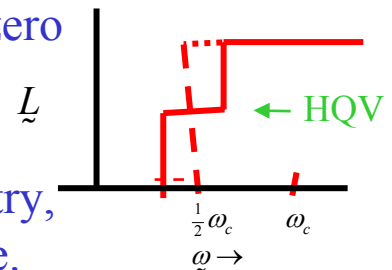
Consider thin annulus rotating at ang. velocity ω , and df. $\omega_c \equiv \hbar / 2mR^2$

At $\omega = \frac{1}{2}\omega_c$ exactly, the nonrotating state and the ordinary "vortex" (p-state) with both spins rotating are degenerate.



But a simple variational argument shows that barring pathology, there exists a nonzero range of ω close to $\frac{1}{2}\omega_c$ where the HQV is more stable than either!

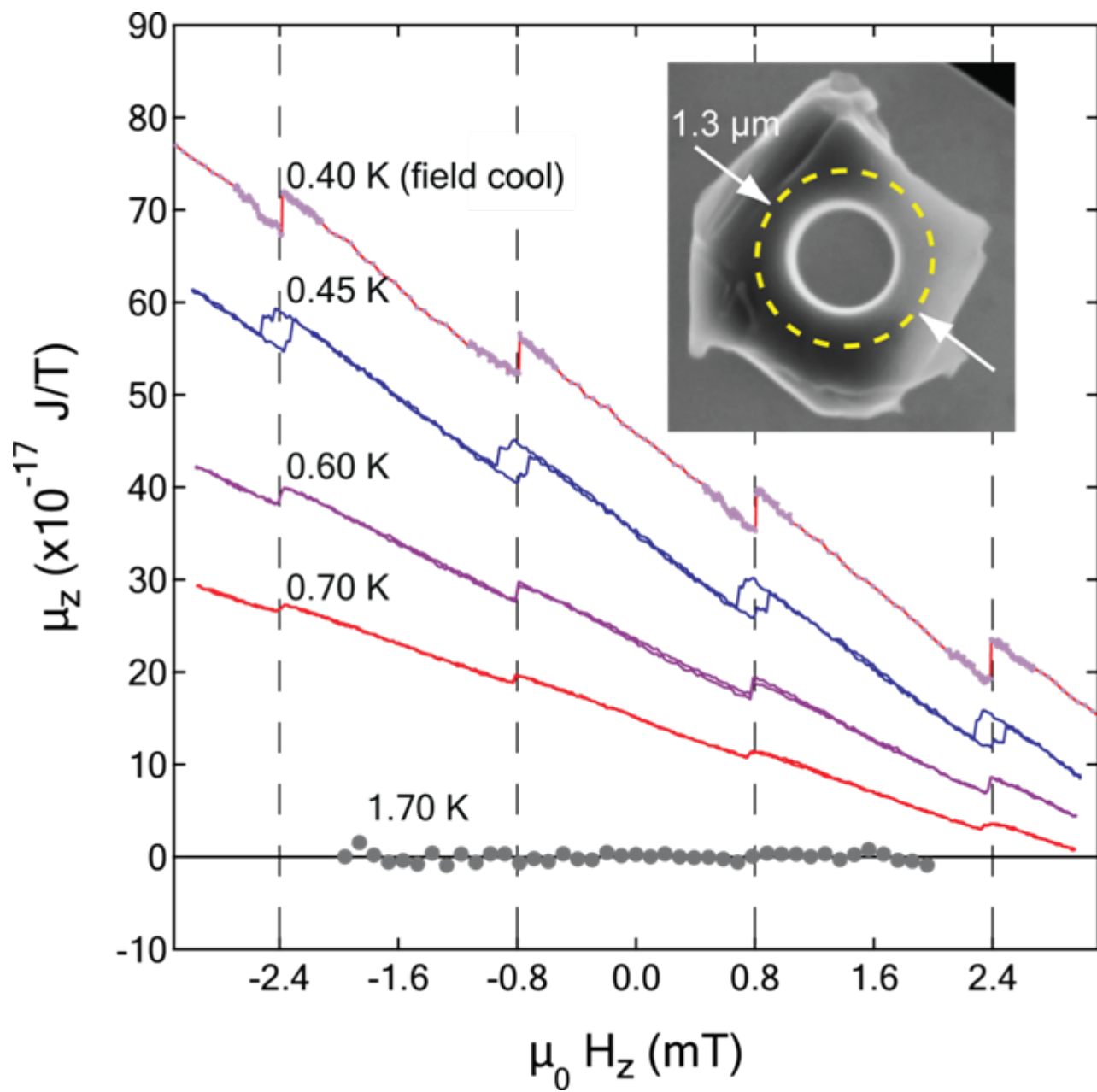
In a simply connected flat-disk geometry, argument is not rigorous but still plausible.



⚠: Yamashita et al. (2008) do experiment in flat-disk geometry, find NO EVIDENCE for HQV!

Possible explanations:

- (1) HQV is not stable under experimental conditions (Kawakami et al., 2009)
- (2) HQV did occur, but NMR detection technique insensitive to it.
- (3) HQV is thermodynamically stable, but inaccessible in experiment.
- (4) Nature does not like HQV's.



QUESTIONS (cont.)

4. Do MF's "exist"?

A. Does existence of a single "Majorana" mode survive replacement of model form $\Delta(\underline{r}, \underline{r}') = \Delta(\underline{R})(\nabla_x + i\nabla_y)\delta(\rho)$ by physical (nonzero-range) form?

B. What is an MF?

(Established approach: 2 MF/s = 1 "real" (Dirac-Bogoliubov) fermion)

Any "completely paired" even-N MBGS can be written

$$\Psi_N \sim \left(\sum_n c_n a_n^+ a_{\bar{n}}^+ \right)^{N/2} |\text{vac}\rangle \quad n + \bar{n} = \text{orthonormal set}$$

Write

$$u_n \equiv (1 + |c_n|^2)^{-1/2}, v_n \equiv c_n (1 + |c_n|^2)^{-1/2}$$

then:

Any operator of form (or linear combination thereof)

$$v_n a_n^+ + u_n a_{\bar{n}} C^\dagger \leftarrow \cong \sum_n c_n a_n^+ a_{\bar{n}}^+$$

identically annihilates GS ("pure annihilator") while any operator of form (or linear combination thereof)

$$u_n^* a_n^+ - v_n a_{\bar{n}} C^\dagger$$

creates (N+1)-particle state ("DB fermion") (but in general **not** energy eigenstate).

The (N+1)-particle energy eigenstates are the particular combinations of DB states that satisfy the BdG equations.



Suppose now we have found a solution of the BdG equations corresponding with energy eigenvalue $E = 0$. What does it represent?

(a) It can create an energy eigenstate of the $N + 1$ -particle system, (DB fermion) with excitation energy 0 relative to the N -particle GS.

(b) It can simply correspond to a **pure annihilator!** But, in neither case can $u(r) = v^*(r)$ as required for an M. F.

Conclusion: An M. F. is nothing but a **quantum superposition of an $E = 0$ DB fermion and a pure annihilator!**

If GS is strongly entangled, DB fermion may be “split” between 2 spatially separated regions (e.g. 2 mutually remote vortices): then MF’s may be localized around single vortices.

So:

$E = 0$ DB fermion + pure annihilator \Rightarrow 2 MF’s.

The \$64K question: **DOES THE CONVERSE HOLD?**

i.e. does existence of 2 MF’s establish existence of single “split” DB fermion?



Can we mimic the real-life 2-vortex problem with an exactly soluble toy model?

Plausible shot: “directed-link” model (Kitaev):



$$\hat{H} = -\sum_j (ta_{j+1}^+ a_j + \Delta a_{j+1}^+ a_j^+) + H.c.$$

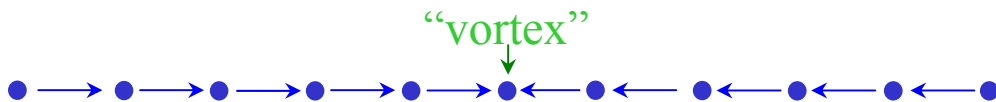
("1D $p + ip$ ")

$$\Delta = (\pm)t$$

Exactly soluble, sustains MF's on ends:



But: obvious way to form “vortex” is to reverse sign of Δ relative to t



Still exactly soluble, but now retains 2 MF's at “vortex”!



If true for real case, disaster for TQC!

? ?

Challenge: find **exactly soluble** model of $(p + ip)$ Fermi superfluid with 2 vortices, and establish existence/nonexistence of single DB fermion “split” between them.