

Superconductivity and Mottness: Exact Results

Nature Physics, vol.16, 1175-1180 (2020);
vol. 18, 511-516 (2022); PRB, 105, 184509.

Luke Yeo



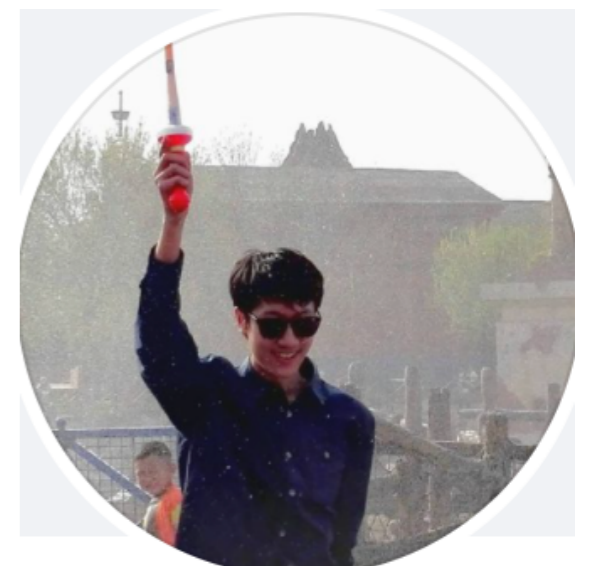
Edwin Huang



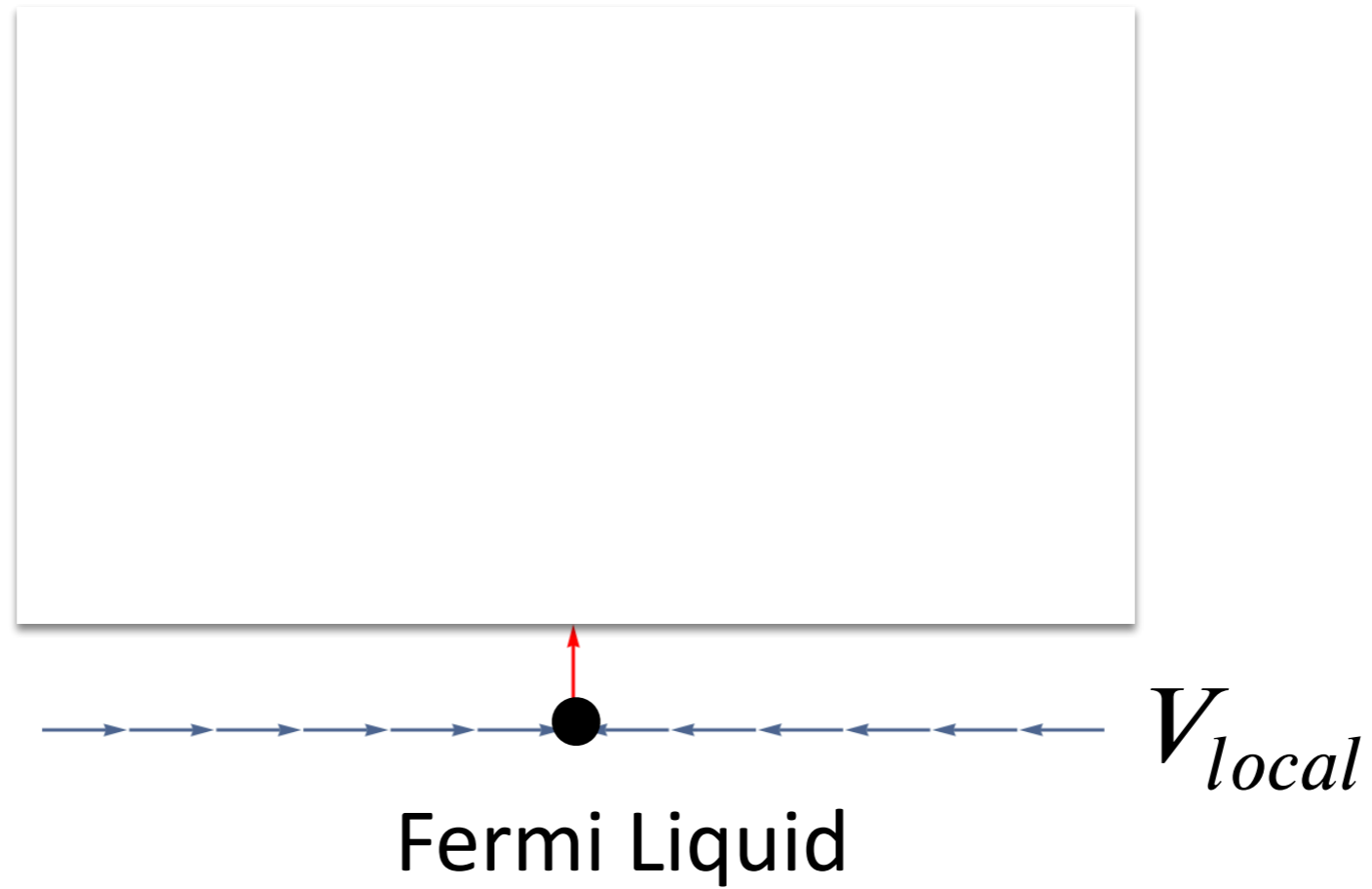
G. La Nave



Jinchao Z.



BCS Superconductor



fixed
point beyond
FL?

quartic
interacting
theory?

Discovery of Orbital Ordering in $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+x}$

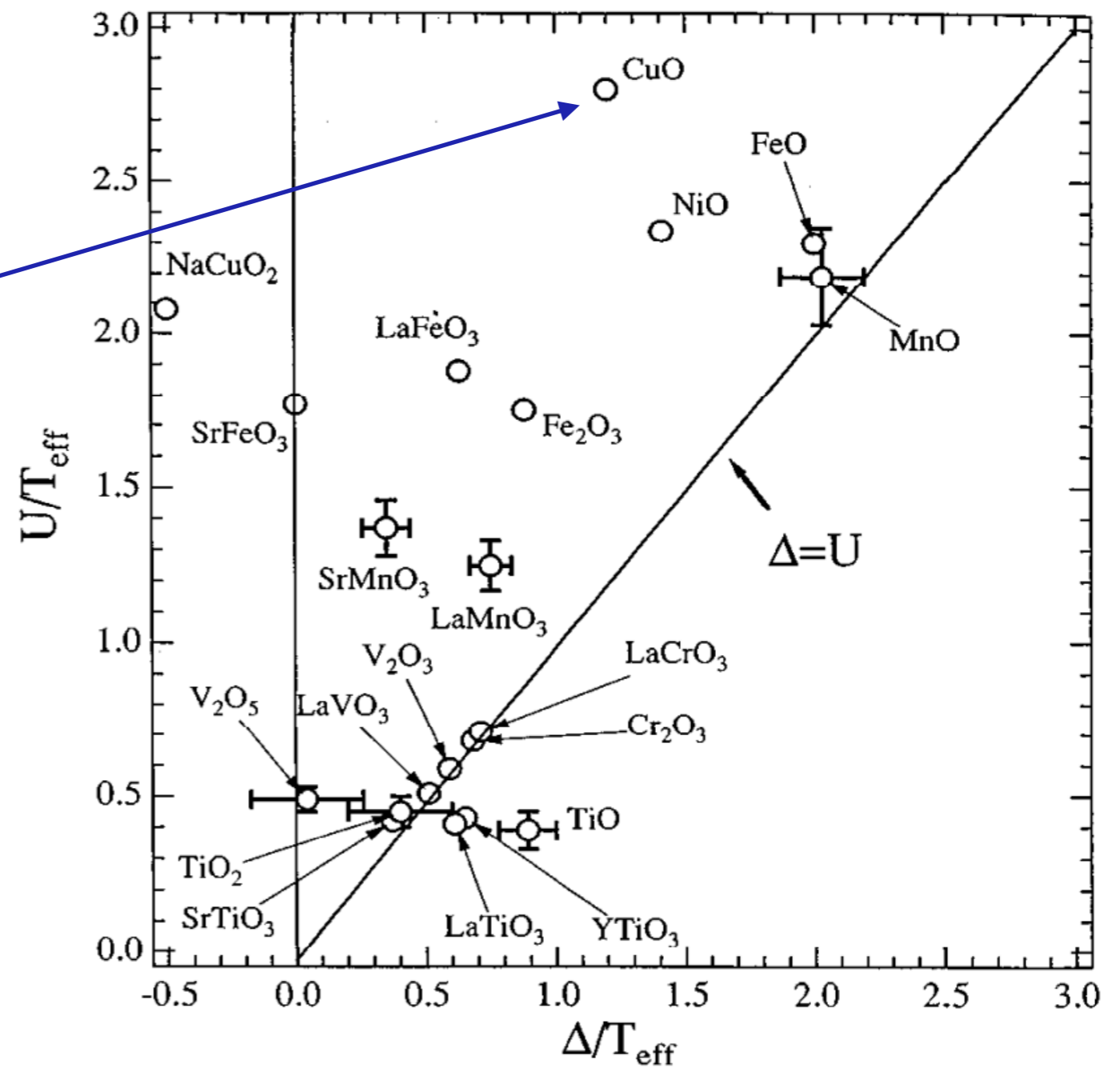
Shuqiu Wang^{1§}, Niall Kennedy^{1,2§}, K. Fujita³, S. Uchida⁴, H. Eisaki⁵,

P.D. Johnson^{1,3}, J.C. Séamus Davis^{1,2,6,7} and S.M. O'Mahony²

Overall, these data reveal a $Q=0$ orbitally ordered state that lifts the energy degeneracy of p_x/p_y oxygen orbitals at separate CuO_2 oxygen sites, in striking analogy to the ordering of d_{zx}/d_{zy} iron orbitals of the iron-based superconductors.

charge-transfer insulator

NFL:
Mottness



solve Hubbard model!!

or instability

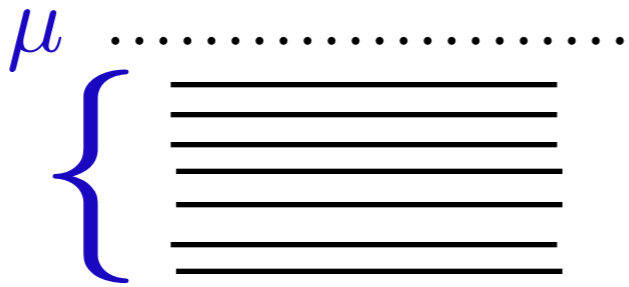
Is there a simpler
approach that has Mottness?

Yes!

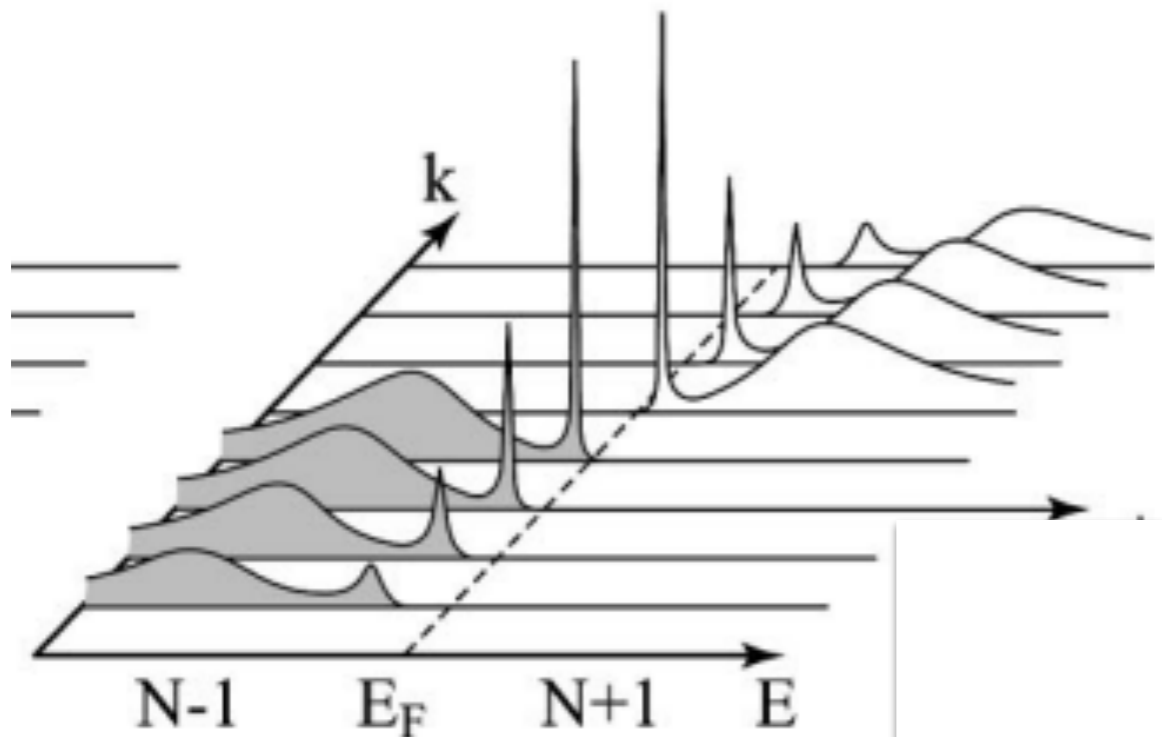
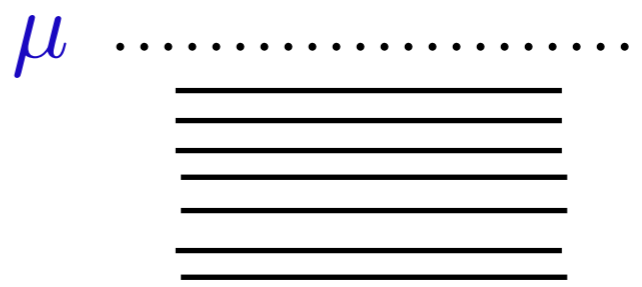
Fermi liquids

NFL

doubly occupied



Is single occupancy below chemical potential possible?

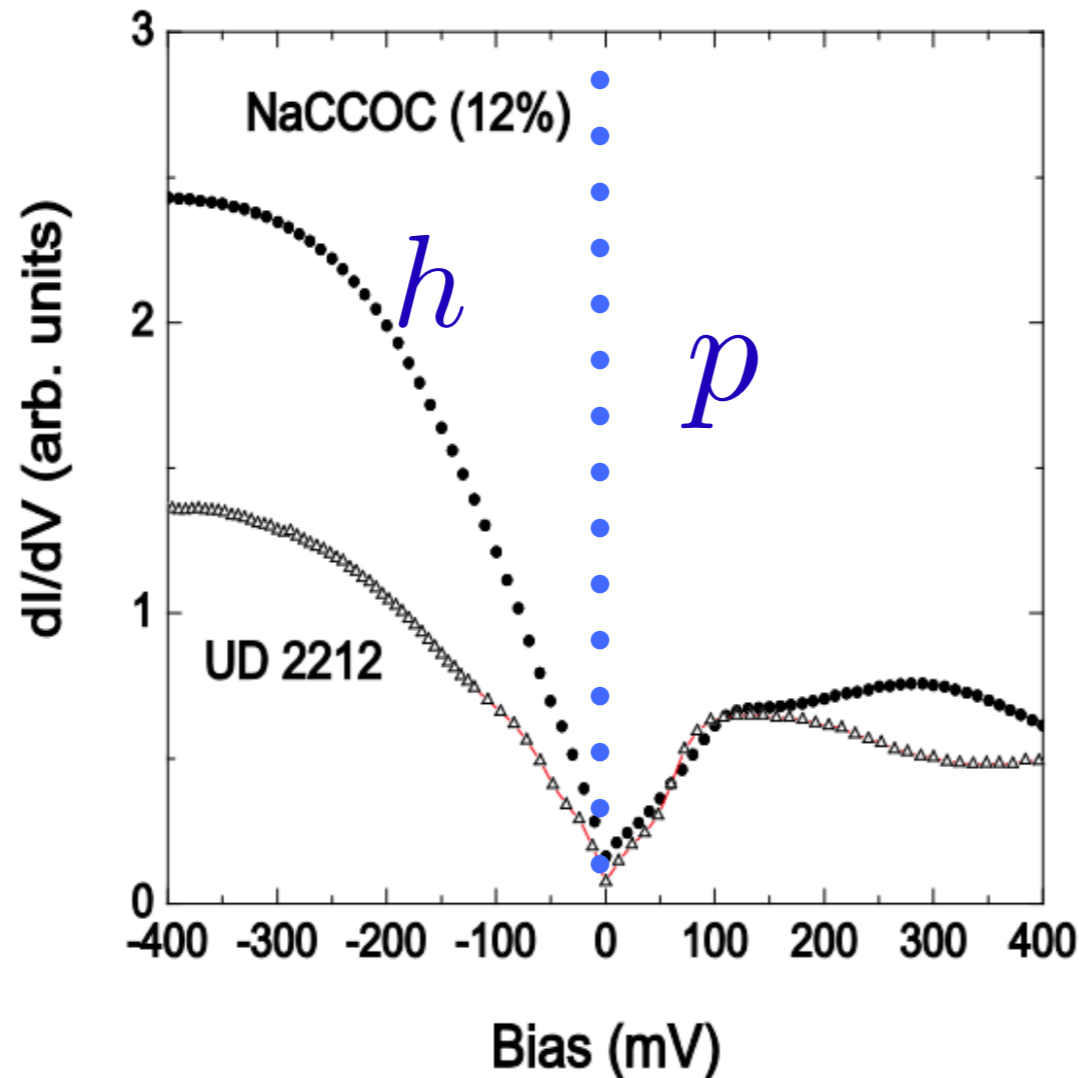


Fermi-liquid system

p-h symmetry

with time-reversal symmetry intact?

p-h asymmetry



LAST WORDS ON THE CUPRATES

P W Anderson, Princeton University

theory. I remain baffled by the almost universal refusal of theorists to confront this evident fact of hole-particle asymmetry head on. Its meaning is that the first step of any

single
occupancy

?

particle-hole
asymmetry



Anderson
Haldane
2000

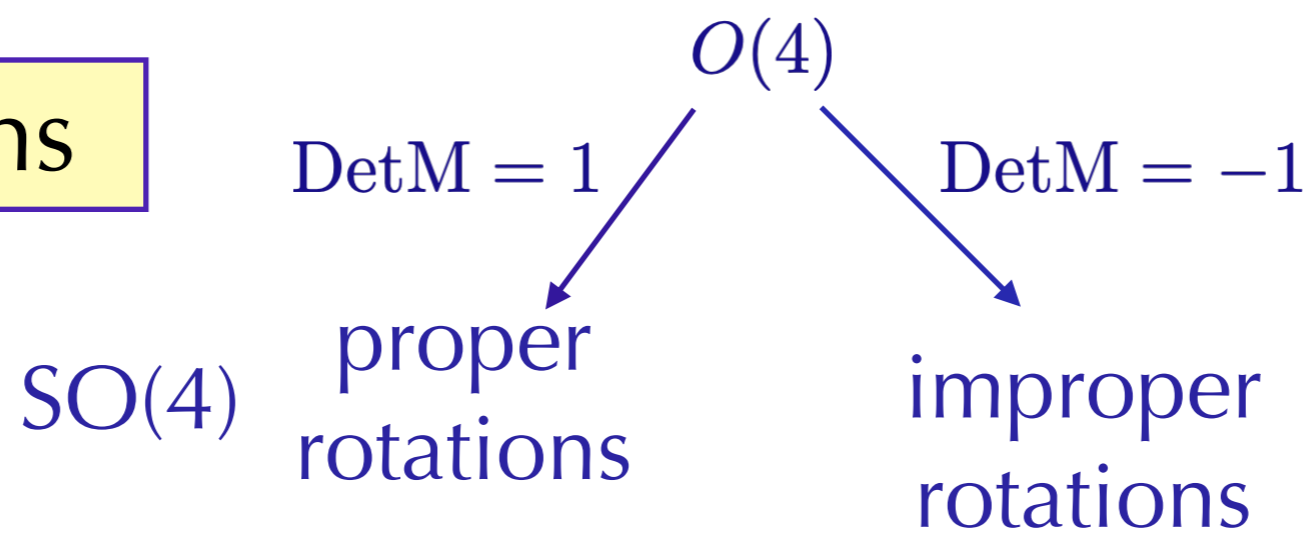
3 citations

Fermi liquids

$$H = \sum_{p,\sigma} (\epsilon(p) - \epsilon_F) n_{p\sigma} + \dots \rightarrow 0$$

$(n_{p\uparrow}, n_{p\downarrow})$ conserved currents

$(c_{p\uparrow}, c_{p\downarrow}, \text{h.c.})$ 4 objects



$$\text{Det}M = \pm 1 \implies Z_2 = O(4) \div SO(4)$$

Improper Rotations

Majorana basis

$$\begin{pmatrix} c_{p\uparrow} \\ c_{p\uparrow}^\dagger \\ c_{p\downarrow} \\ c_{p\downarrow}^\dagger \end{pmatrix} \longrightarrow \begin{pmatrix} c_{p\uparrow} + c_{p\uparrow}^\dagger \\ i(c_{p\uparrow} - c_{p\uparrow}^\dagger) \\ c_{p\downarrow} + c_{p\downarrow}^\dagger \\ i(c_{p\downarrow} - c_{p\downarrow}^\dagger) \end{pmatrix}$$

$$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} c_{p\uparrow} + c_{p\uparrow}^\dagger \\ i(c_{p\uparrow} - c_{p\uparrow}^\dagger) \\ c_{p\downarrow} + c_{p\downarrow}^\dagger \\ i(c_{p\downarrow} - c_{p\downarrow}^\dagger) \end{pmatrix} \longrightarrow c_{p\downarrow} \rightarrow c_{p\downarrow}^\dagger$$

p-h transformation

$$\epsilon(p) = \epsilon_F$$

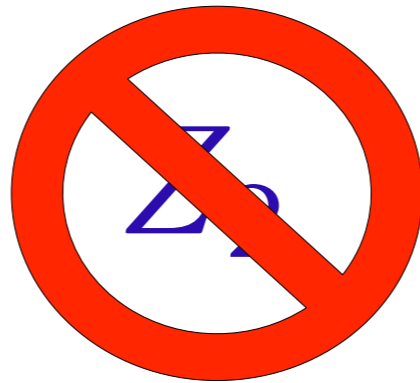
Fermi
Surface

$$H = 0$$



$$\left. \begin{array}{l} n_{p\uparrow} \rightarrow 1 - n_{p\uparrow} \\ n_{p\downarrow} \rightarrow n_{p\downarrow} \end{array} \right\} \mathbb{Z}_2 \text{ at Fermi surface only}$$

How to destroy Fermi liquids?



$$H = \sum_{p,\sigma} (\epsilon(p) - \epsilon_F) n_{p\sigma} + U n_{p\uparrow} n_{p\downarrow}$$

odd
under Z_2

scaling dimension

$$[n_{p\uparrow} n_{p\downarrow}] = -2$$

relevant
interaction

New fixed point!

Hatsugai-Kohmoto
model

Hubbard
not
necessary!

General HK Model

$$\sum_k (\xi_k (n_{k\uparrow} + n_{k\downarrow}) + U n_{k\uparrow} n_{k\downarrow})$$

Solvable Mott transition: $U > W$

$$G_{k\sigma}(i\omega_n \rightarrow z) = \frac{1 - \langle n_{k\bar{\sigma}} \rangle}{z - \xi_k} + \frac{\langle n_{k\bar{\sigma}} \rangle}{z - (\xi_k + U)} \neq \frac{1}{z - \omega_k}$$

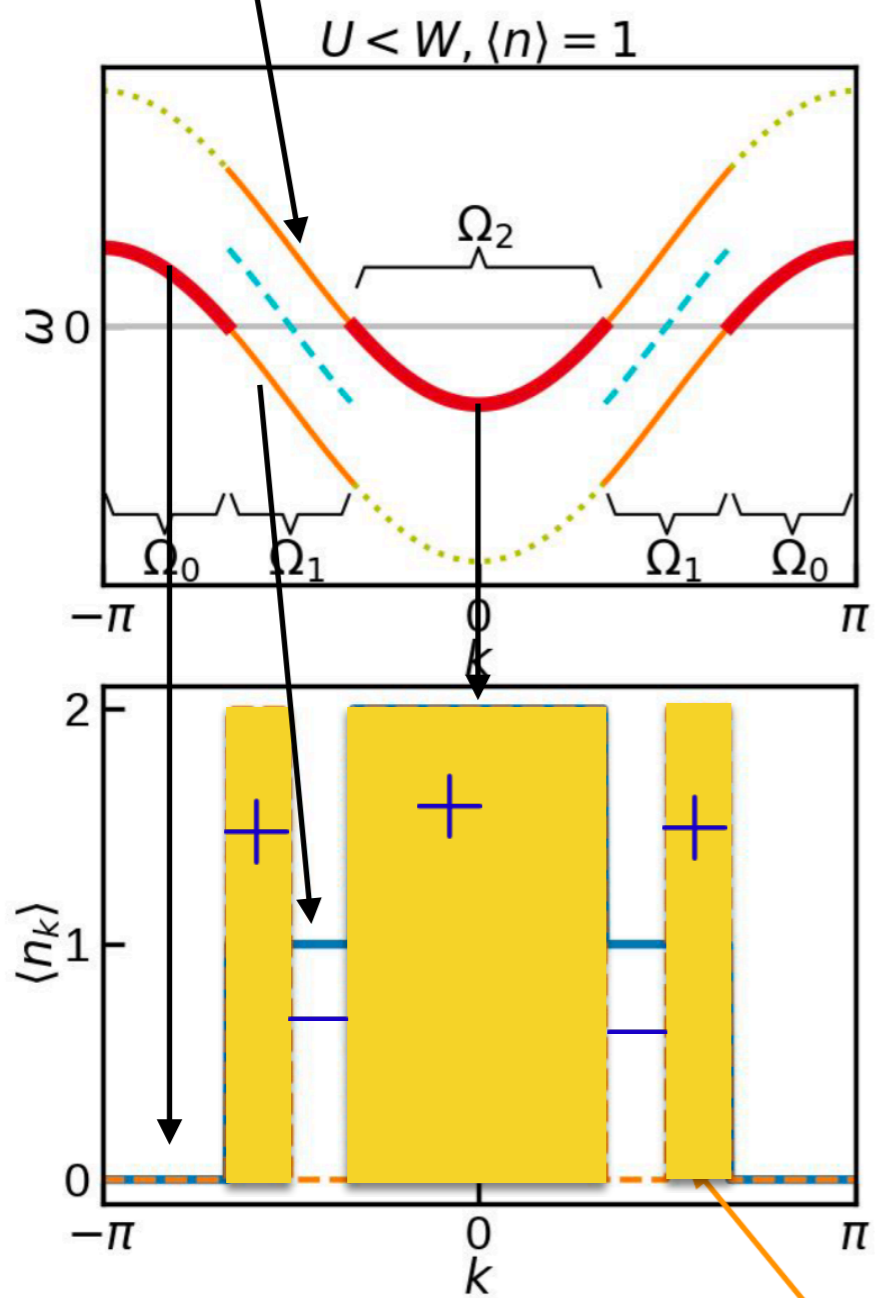
lower Hubbard band

upper Hubbard band

zeros

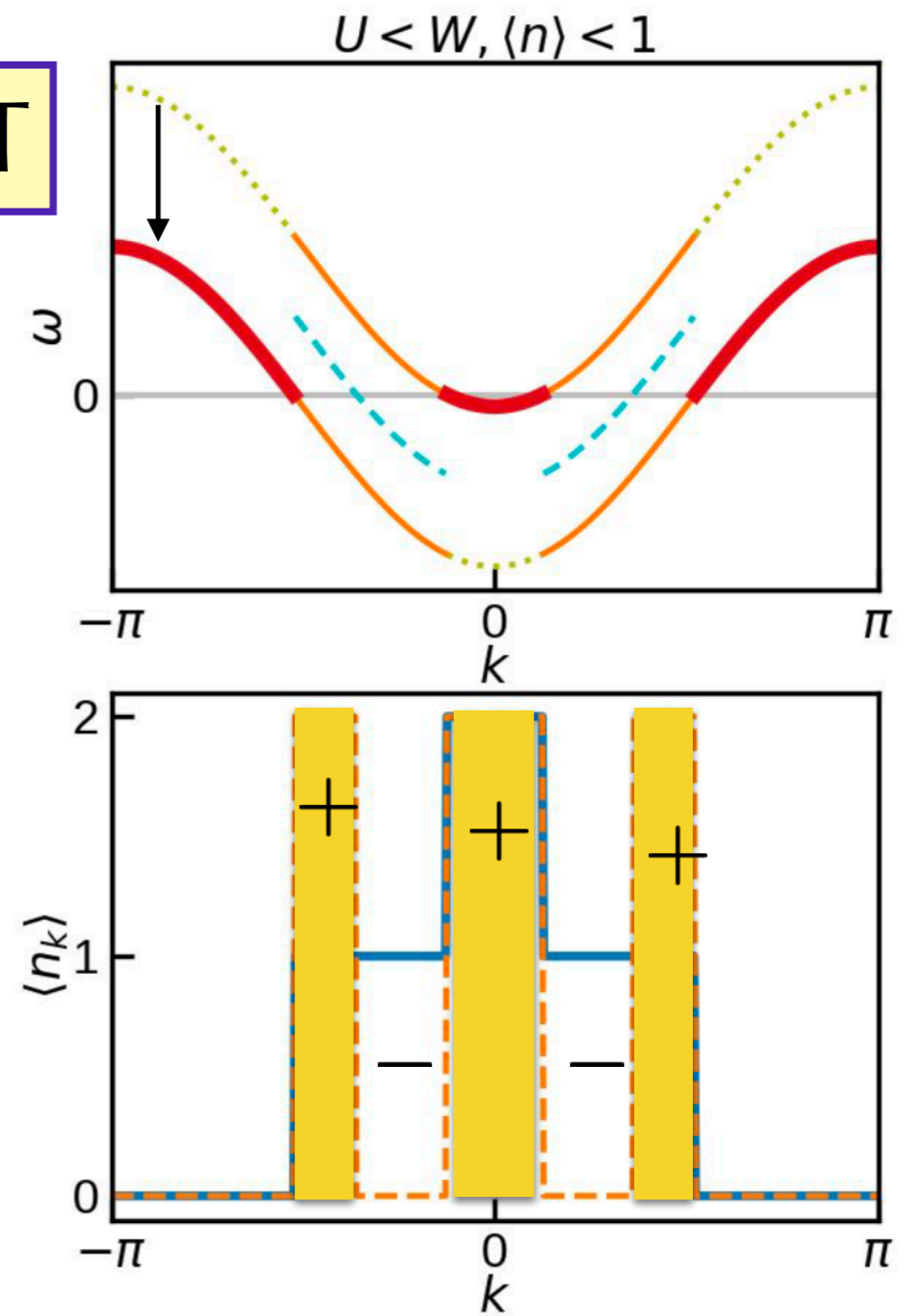
single occupancy

counting charges



$$n_{\text{Lutt}} = \langle n \rangle$$

SWT



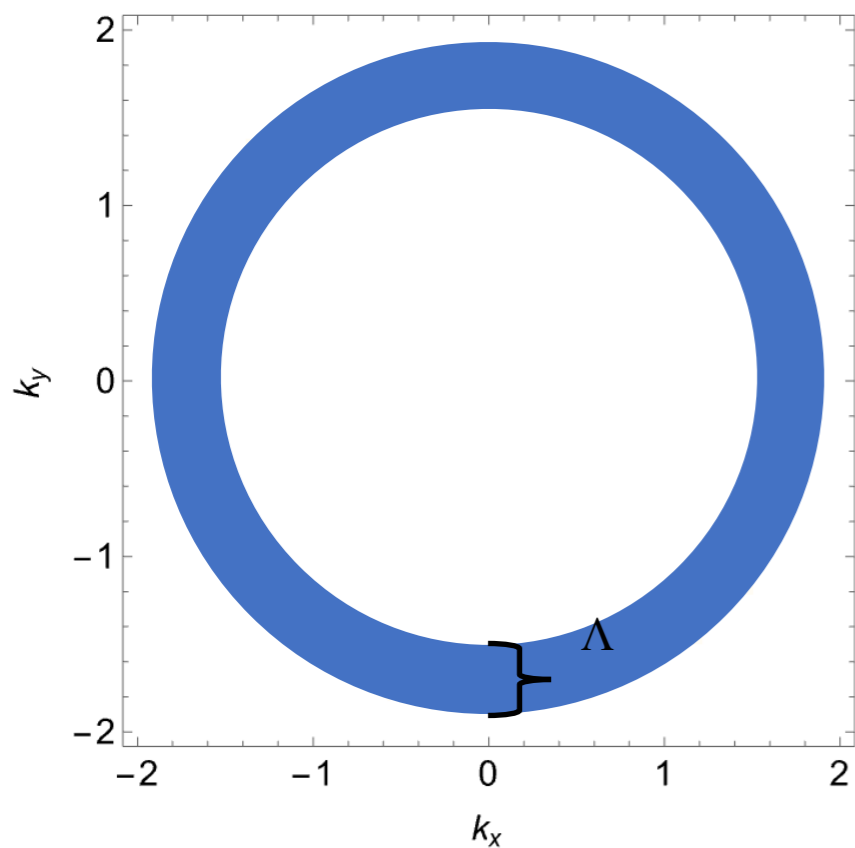
$$n_{\text{Lutt}} \neq \langle n \rangle$$

zeros \neq particles

$$H_{\text{HK}} + V_{\text{local}}$$

$$\delta S_4 = \frac{1}{4} \int u(4, 3, 2, 1) \bar{\psi}(4) \bar{\psi}(3) \psi(2) \psi(1) \times$$

$$\prod_{i=1}^4 \frac{d^d k_i d\omega_i}{(2\pi)^{d+1}} \delta(\mathbf{k}_1 + \mathbf{k}_2 - \mathbf{k}_3 - \mathbf{k}_4) \delta(\omega_1 + \omega_2 - \omega_3 - \omega_4)$$



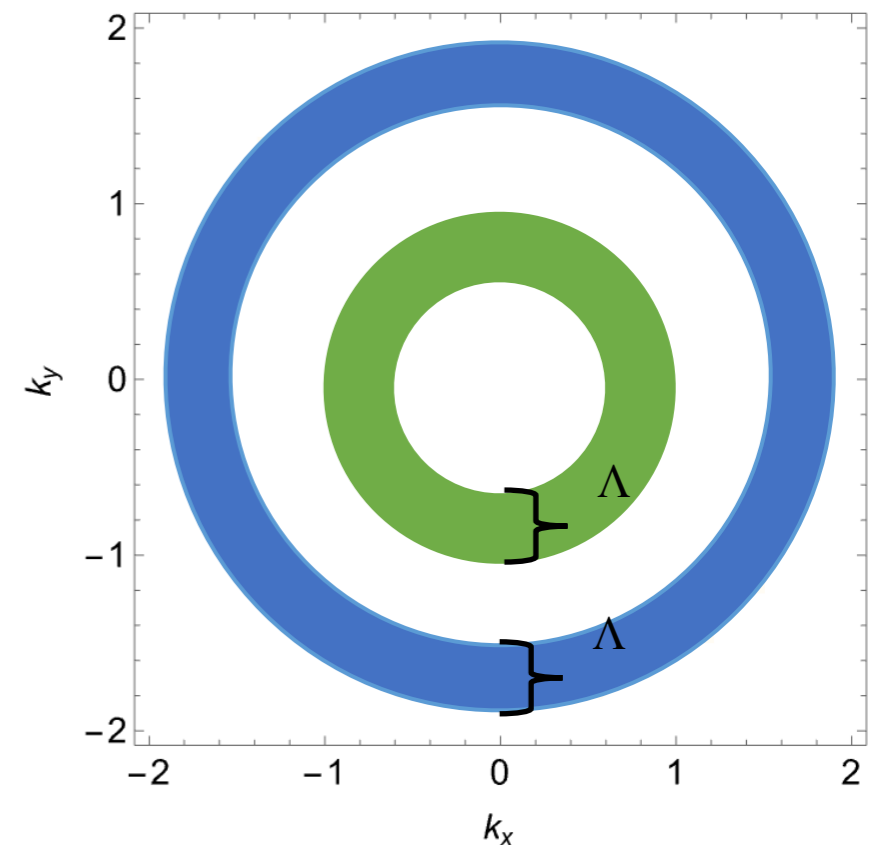
$$\Lambda \rightarrow \Lambda/s$$

$$q' = qs$$

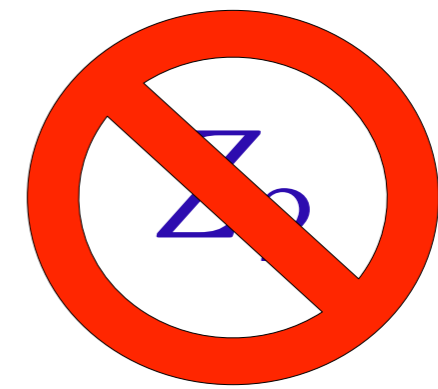
$$\omega' = \omega/s$$

$$\psi'(k', \omega') = s^{-3/2} \psi(k, \omega)$$

$$U' = s^2 U$$



Local (Hubbard) part of 4-Fermion interactions are still irrelevant



FL

HK

quantum numbers

k

k

stability

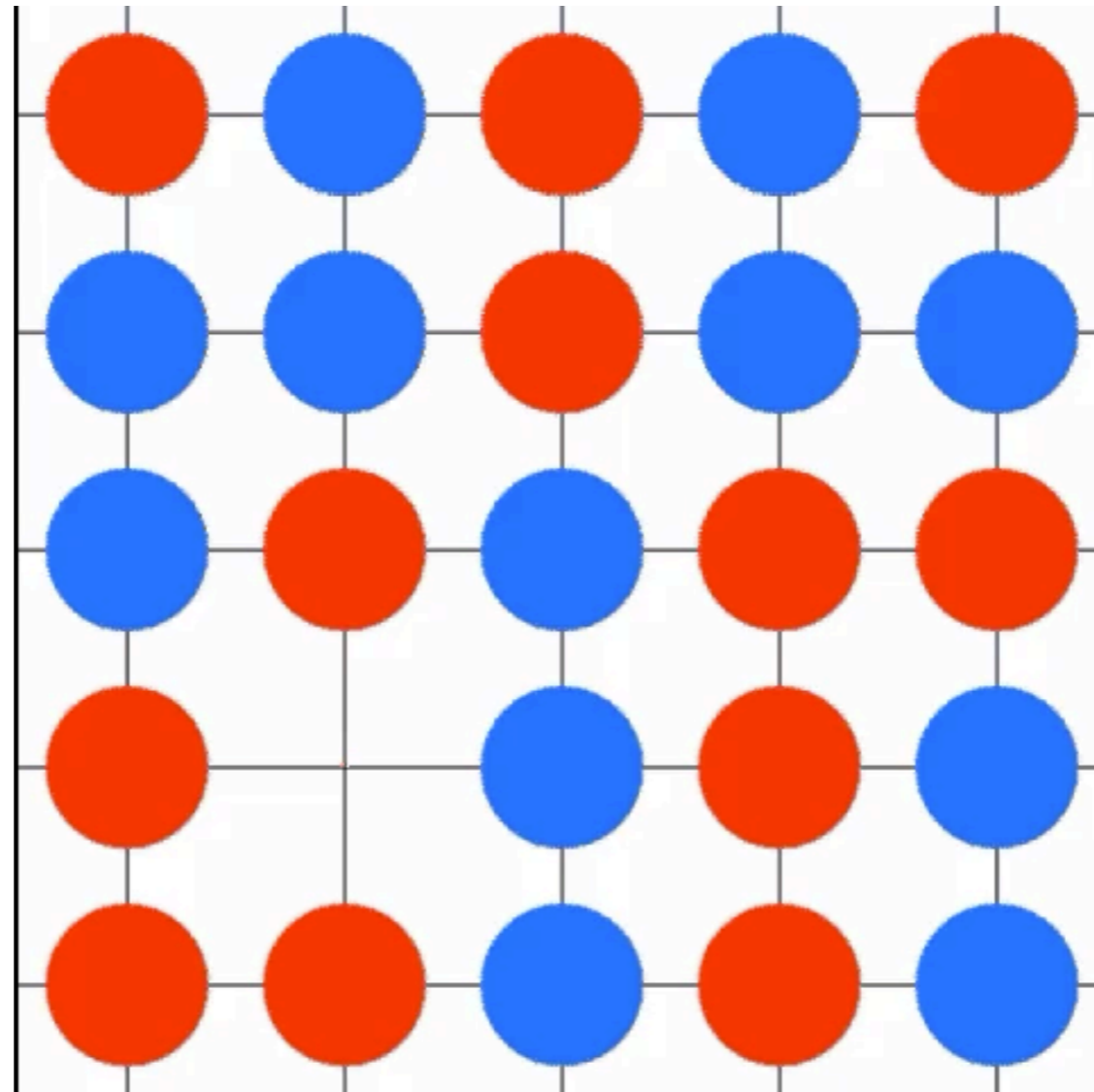
$$\left\{ \begin{array}{l} V_{\text{local}} \quad \beta(V_{\text{local}}) = 0 \\ V_{\text{sc}} \quad \beta(V_{\text{sc}}) \rightarrow \infty \end{array} \right.$$

$$\beta(V_{\text{local}}) = 0$$

?

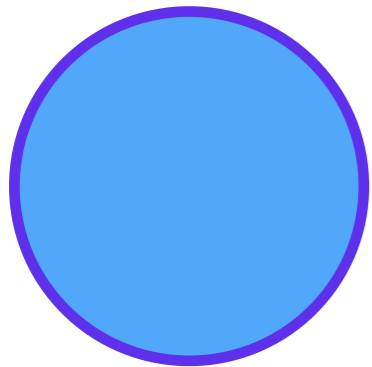
what does the HK model leave out??

$$[H_t, H_U] \neq 0$$

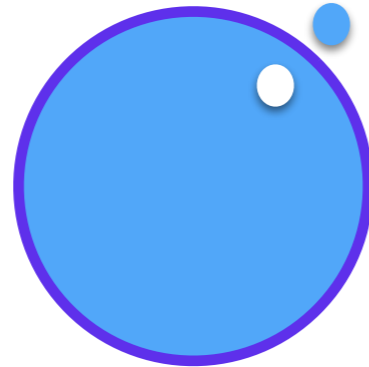


dynamical spectral weight transfer

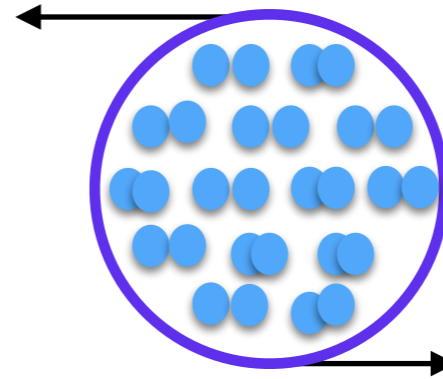
Fermi gas



Fermi liquid



BCS
superconductor



Mottness

2

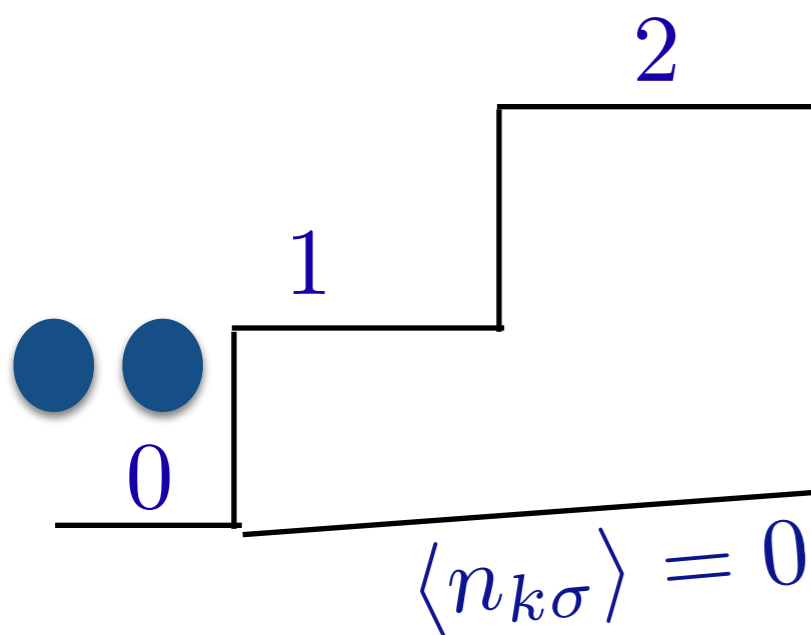
1

0



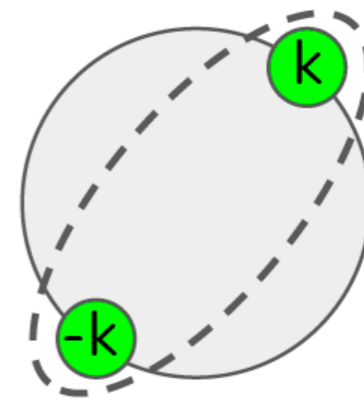
Superconductivity?

Cooper Instability



$$H = H_{\text{HK}} - gH_p$$

$$|\psi\rangle = \sum_{k \in \Omega_0} \alpha_k b_k^\dagger |\text{GS}\rangle$$

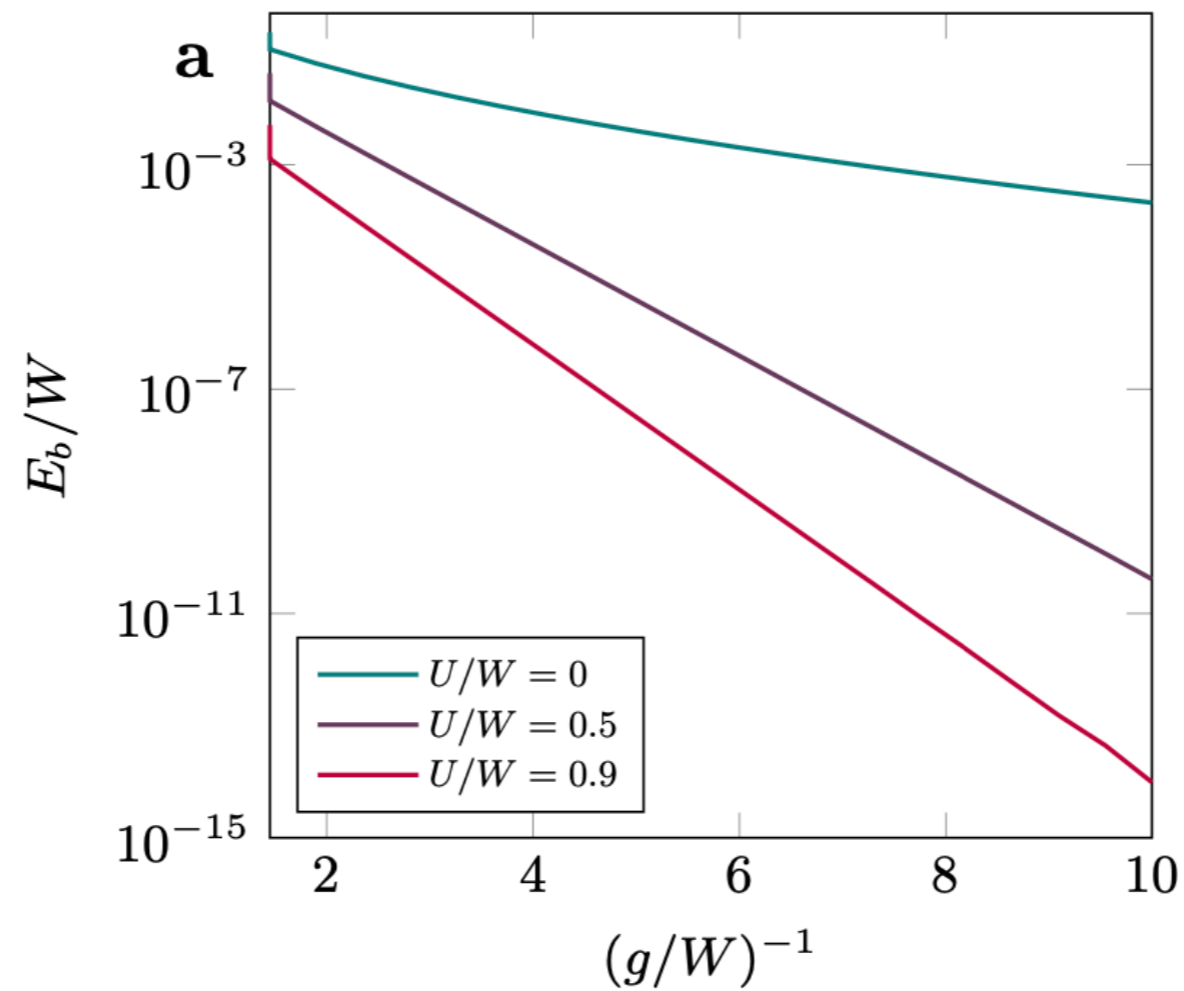


$$= \sum_{k, k'} b_k^\dagger b_{k'}$$

$$E_b = \langle \text{GS} | H | \text{GS} \rangle - \langle \psi | H | \psi \rangle \leq 0$$

Cooper Instability

$$E_b = -E \sim W(1 - (U/W)^2)e^{-\pi W \sqrt{1 - (U/W)^2} / g}$$



Pair Susceptibility

$$\chi(i\nu_n) \equiv \frac{1}{L^d} \int_0^\beta d\tau e^{i\nu_n \tau} \langle T \Delta(\tau) \Delta^\dagger \rangle_g$$

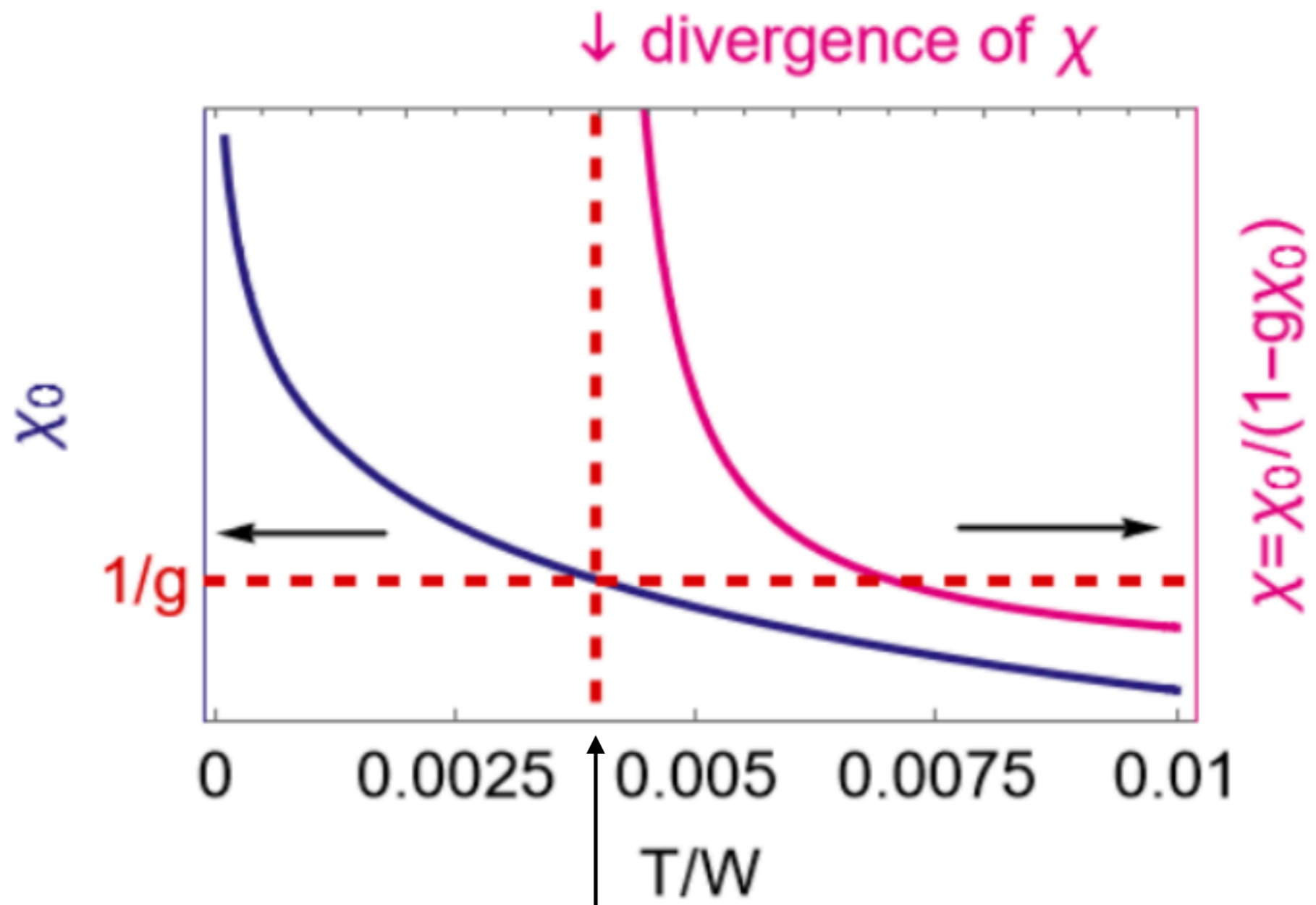


$$= \frac{\chi_0}{1 - g\chi_0}$$



$$g\chi_0 = 1$$

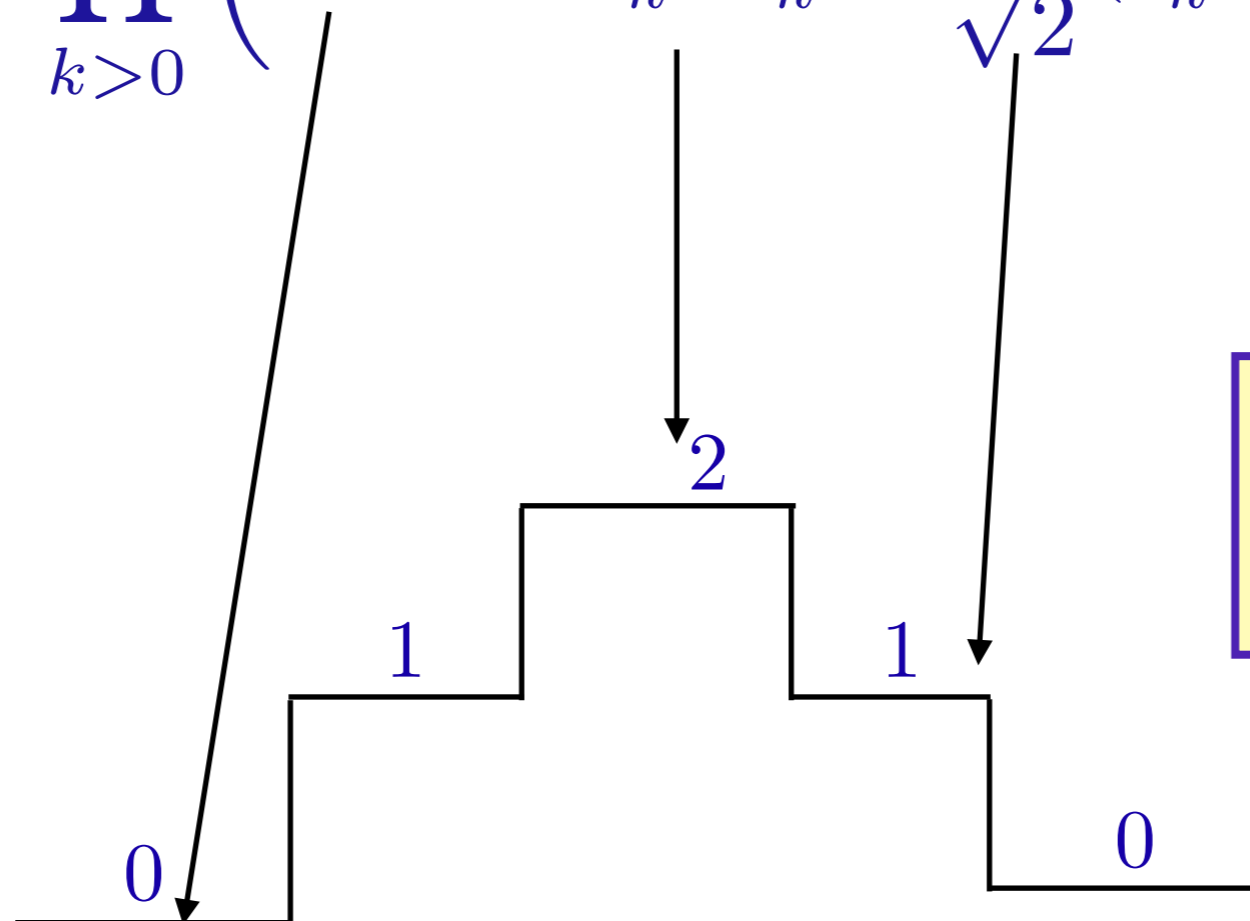
solve for T_c



$$T_c = (W - U)^{4/5} U^{1/5} \frac{e^\gamma}{\pi} e^{-\frac{4}{5} \frac{W}{g}}.$$

variational MF wave function

$$|\psi\rangle = \prod_{k>0} \left(x_k + y_k b_k^\dagger b_{-k}^\dagger + \frac{z_k}{\sqrt{2}} (b_k^\dagger + b_{-k}^\dagger) \right) |0\rangle$$



HK
generalization

three variational parameters

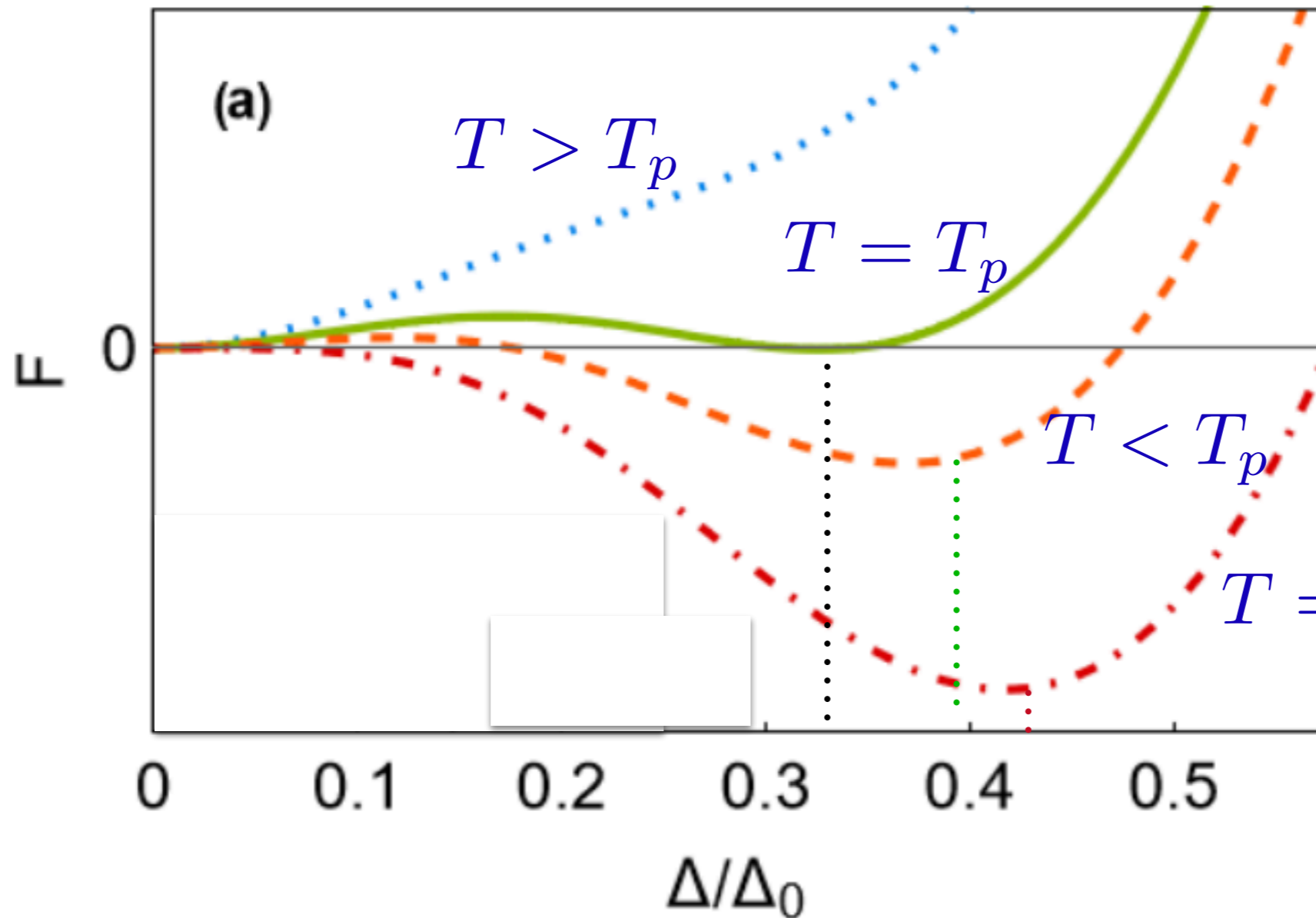
$$|x_k|^2 + |y_k|^2 + |z_k|^2 = 1$$

gap equation


$$\Delta \ll U, W$$

$$\Delta = (W - U)^{1/2} U^{1/2} e^{-\frac{W}{2g}}$$

compute free energy



$t_G \approx 10^{-11}$

MF theory
is accurate!

$$\lim_{g \rightarrow 0} \frac{\Delta}{T_c} \rightarrow \infty$$

Two-stage superconductivity in the Hatsugai-Kohomoto-BCS model

Yu Li,¹ Vivek Mishra,¹ Yi Zhou,^{2,3,4} and Fu-Chun Zhang^{1,4,*}

¹Kavli Institute for Theoretical Sciences, University of Chinese Academy of Sciences, Beijing 100190, China

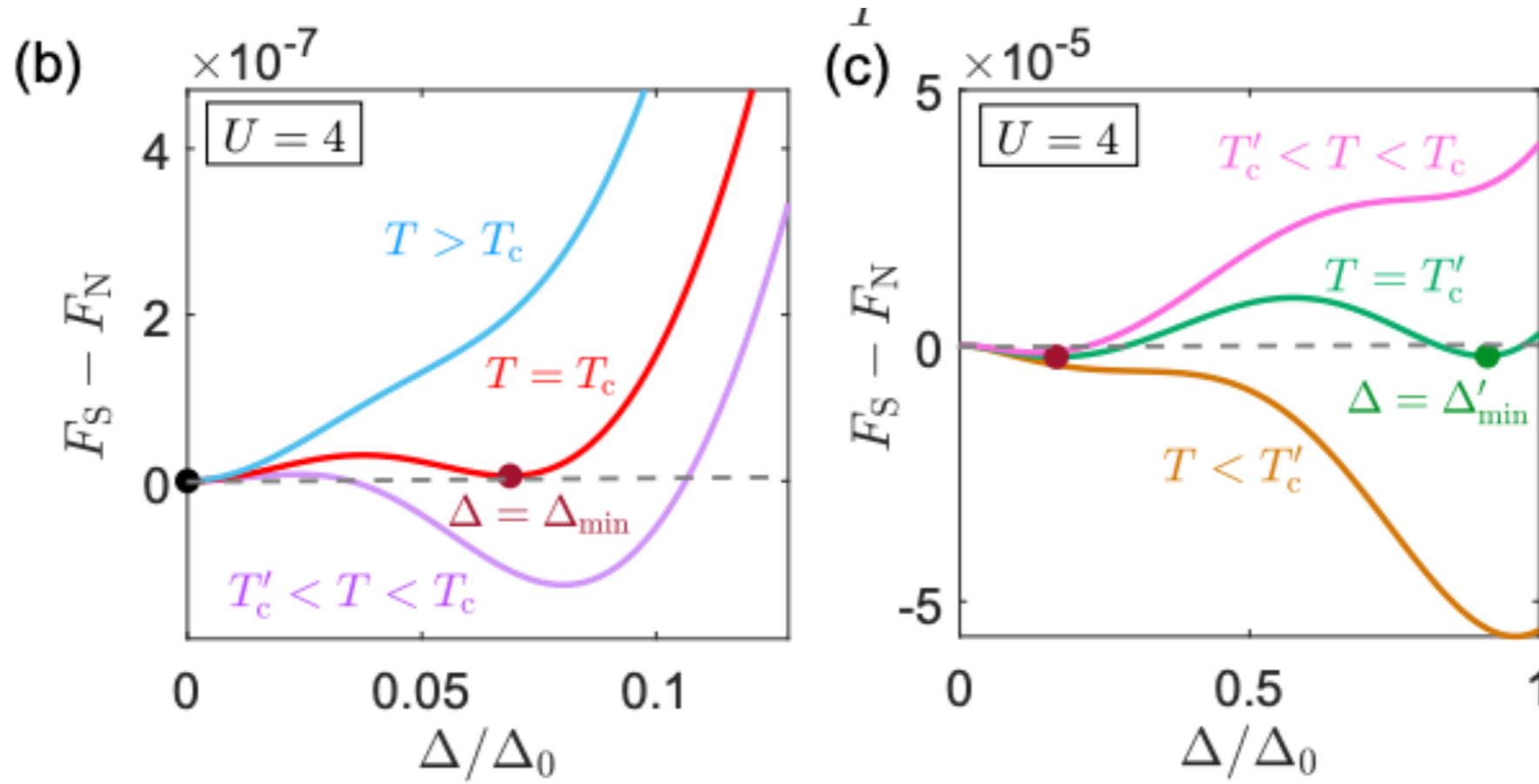
²Beijing National Laboratory for Condensed Matter Physics & Institute of Physics, Chinese Academy of Sciences, Beijing 100190, China

³Songshan Lake Materials Laboratory, Dongguan, Guangdong 523808, China

⁴CAS Center for Excellence in Topological Quantum Computation,
University of Chinese Academy of Sciences, Beijing 100190, China

(Dated: July 7, 2022)

<https://arxiv.org/pdf/2207.01904.pdf>



Bogoliubov excitations

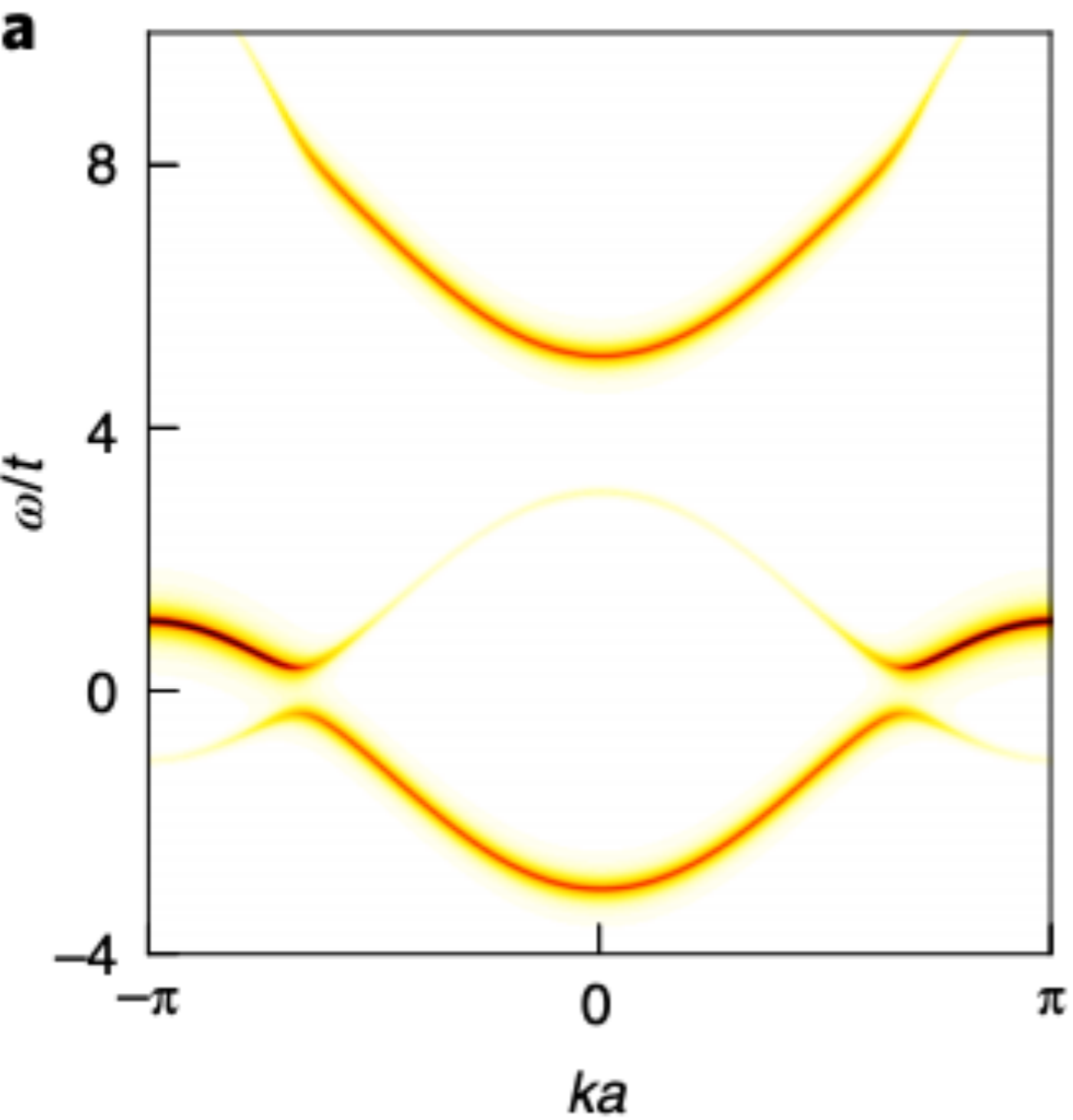
$$\gamma_{k\sigma} |\psi_{\text{BCS}}\rangle = 0$$

$$\gamma_{k\sigma} = u_k c_{k\sigma} - \sigma v_k c_{-k\bar{\sigma}}^\dagger$$


PYHons excitations

$$\gamma_{k\sigma}^l \propto \sqrt{2} x_k \zeta_{k\sigma}^\dagger - \sigma z_k \zeta_{-k\bar{\sigma}}$$

$$\gamma_{k\sigma}^u \propto z_k \eta_{k\sigma}^\dagger - \sigma \sqrt{2} y_k \eta_{-k\bar{\sigma}}$$



PYHon band

can we explain the color change?

REPORT

Superconductivity-Induced Transfer of In-Plane Spectral Weight in $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$

H. J. A. Molegraaf¹, C. Presura¹, D. van der Marel^{1,*}, P. H. Kes², M. Li²

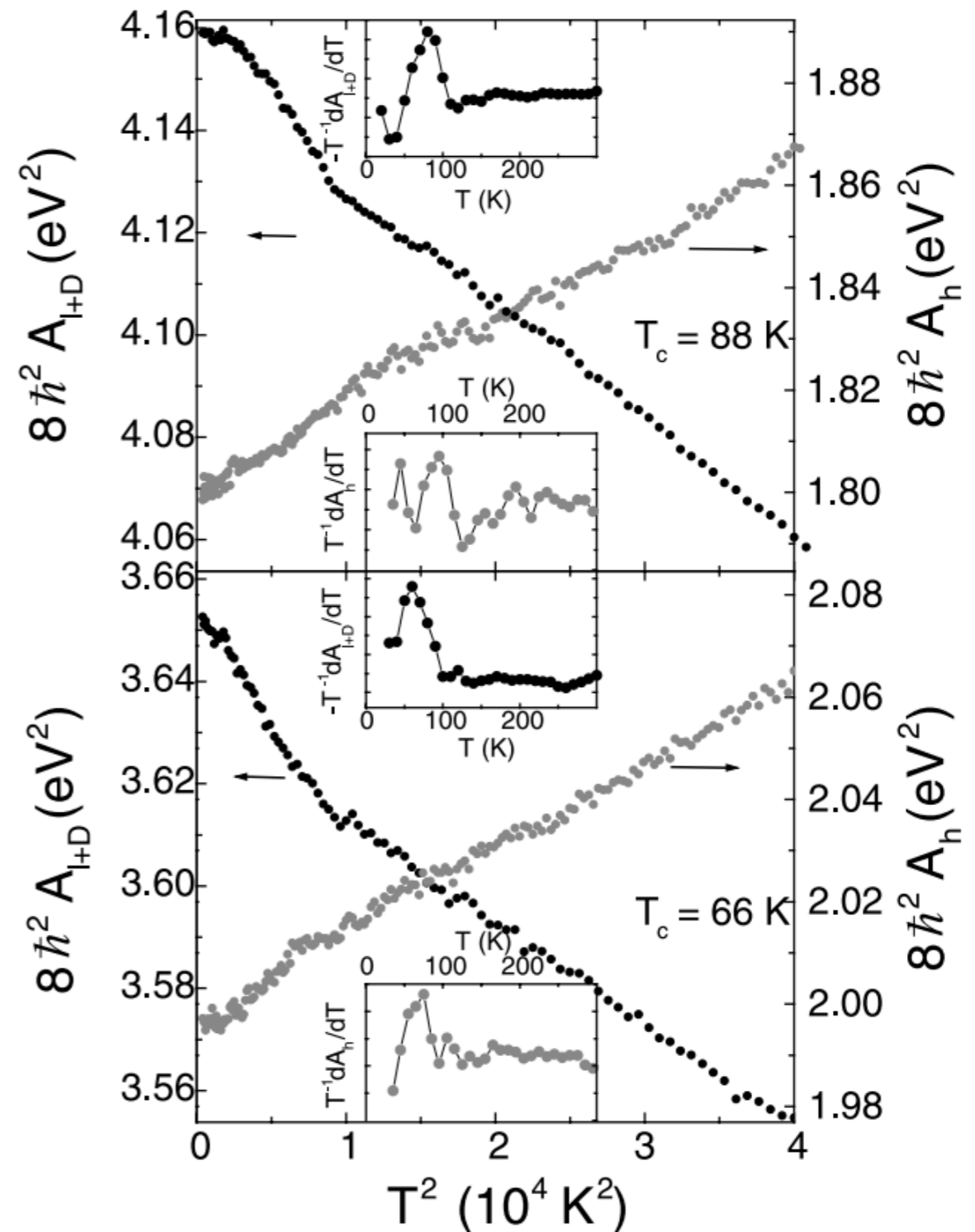
+ See all authors and affiliations

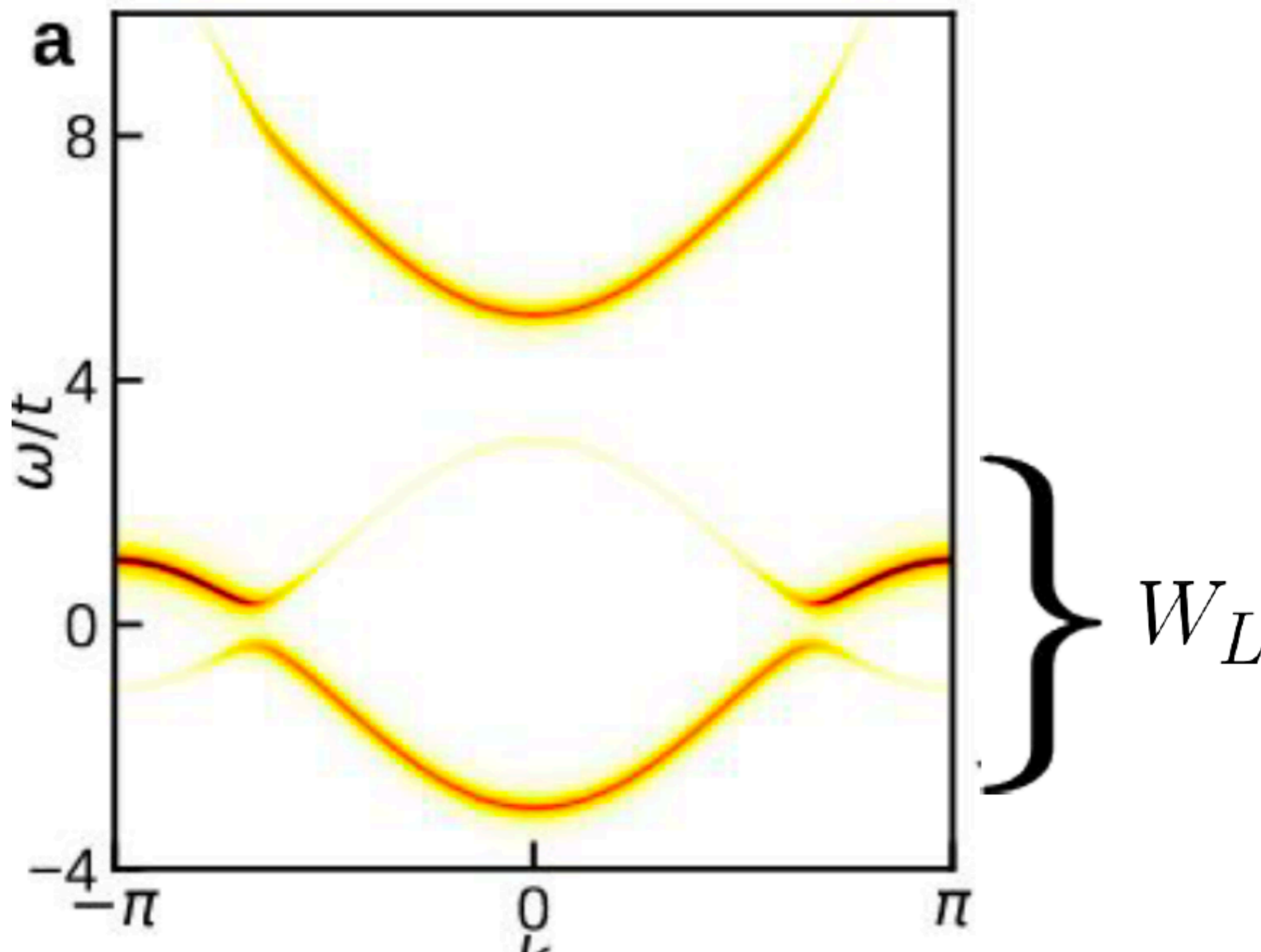
Science 22 Mar 2002:
Vol. 295, Issue 5563, pp. 2239-2241
DOI: 10.1126/science.1069947

$$A_l = \int_0^{\Omega} \sigma(\omega) d\omega \quad \Omega/2\pi c = 10000 \text{ cm}^{-1}$$

$$A_h = \int_{\Omega}^{2\Omega} \sigma(\omega) d\omega \quad \Omega/2\pi c = 10000 \text{ cm}^{-1}$$

$$\frac{\Delta A_l}{A_l} \propto 3\%$$





why?

$$H = H_{\text{HK}} + H_p$$

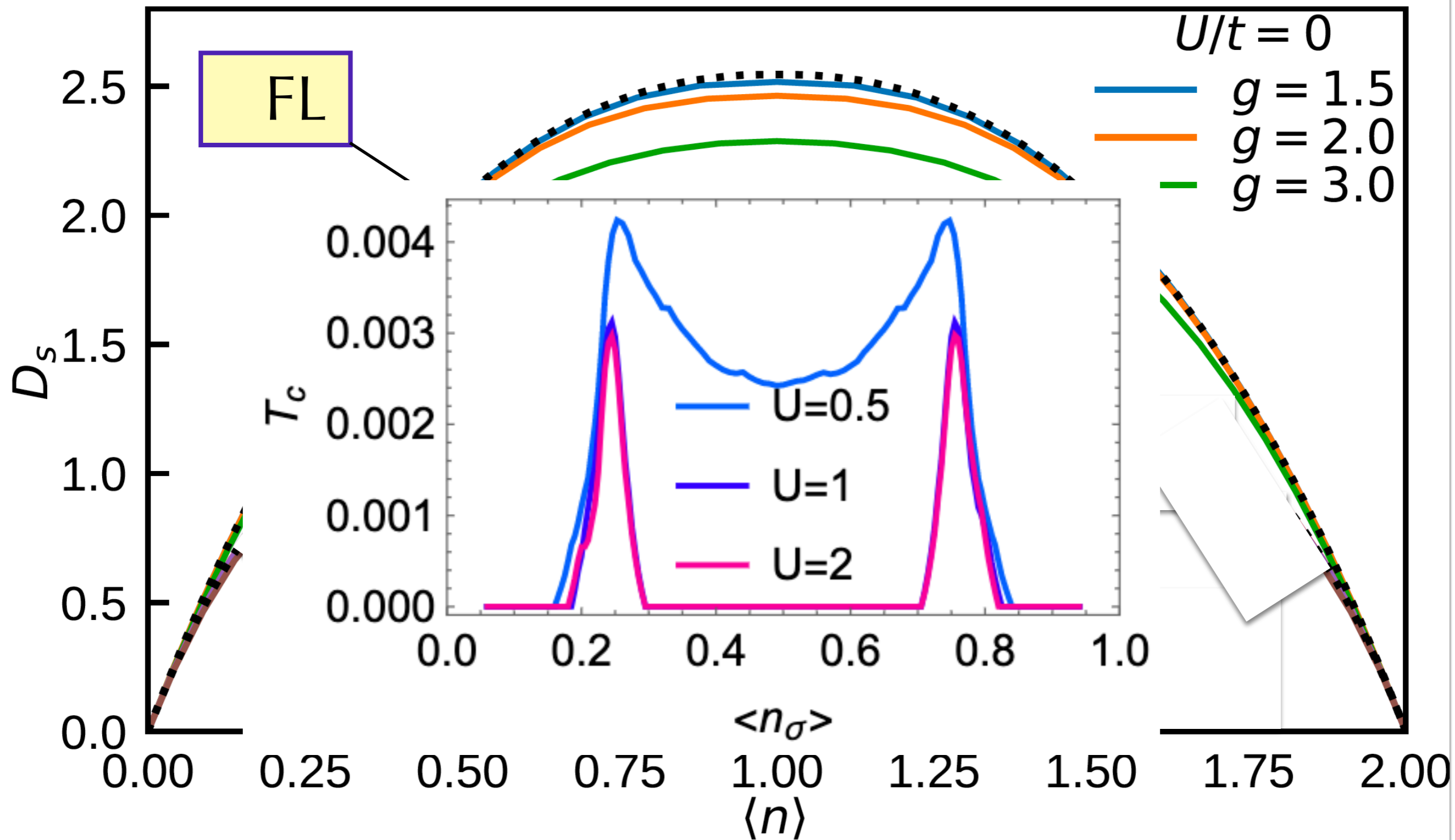
$$[H_{\text{HK}}, H_p] \neq 0$$



dynamical
spectral weight
transfer

Superfluid Density

Mottness-induced suppression



Superconductivity

Mottness

observable

$$\chi \rightarrow \infty$$

$$\Delta \neq 0$$

$$\lim_{g \rightarrow 0} 2\Delta_0/k_B T_c$$

quasi – particles

t_G (Ginzburg)

$$1/TT_1$$

Landau Expansion

$$E_{\text{cond}}/N(0)\Delta^2$$

BCS/FL

$$T_c$$

$$T_c$$

$$3.52$$

Bogoliubons

$$\approx 10^{-12}$$

HS peak

$$a = \alpha t, b > 0$$

$$-1$$

PYHZ/HK

$$T_c (= T_2)$$

$$T_p (> T_2)$$

$$\infty$$

PYHons

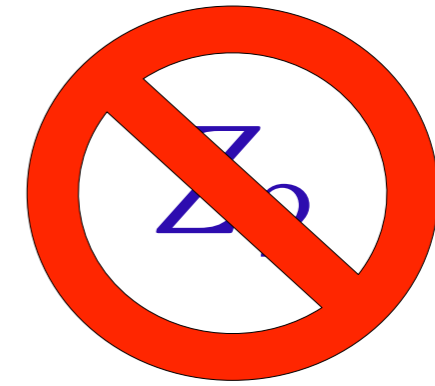
$$\approx 10^{-11}$$

no HS peak

$$a = \alpha t, b < 0$$

$$c > 0$$

$$[-2, -1]$$



FL

HK

quantum numbers

k

k

stability $\left\{ \begin{array}{l} V_{\text{local}} \quad \beta(V_{\text{local}}) = 0 \\ V_{\text{sc}} \quad \beta(V_{\text{sc}}) \rightarrow \infty \end{array} \right.$

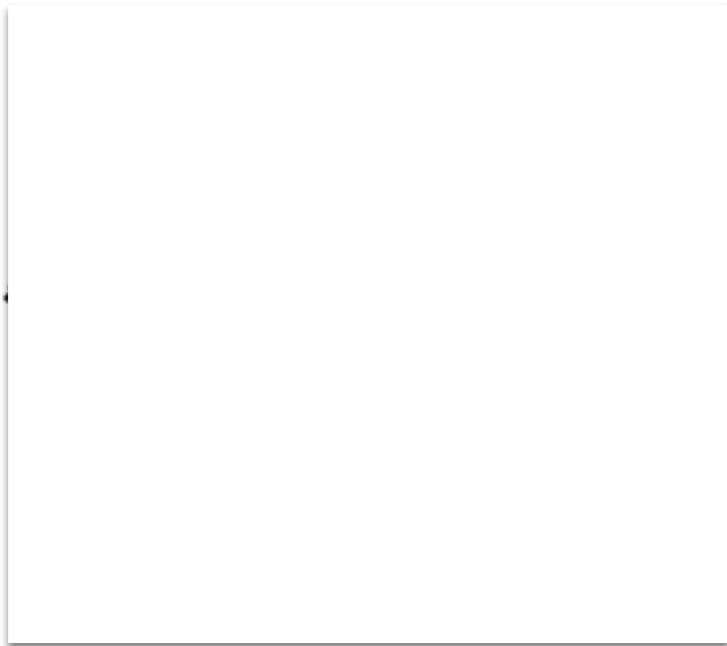
$\beta(V_{\text{local}}) = 0$

$\beta(V_{\text{sc}}?) \rightarrow \infty$

Topology + Strong Correlations?

Are Exact Statements Possible?

Haldane + HK model



TI

1/4-filled MI

